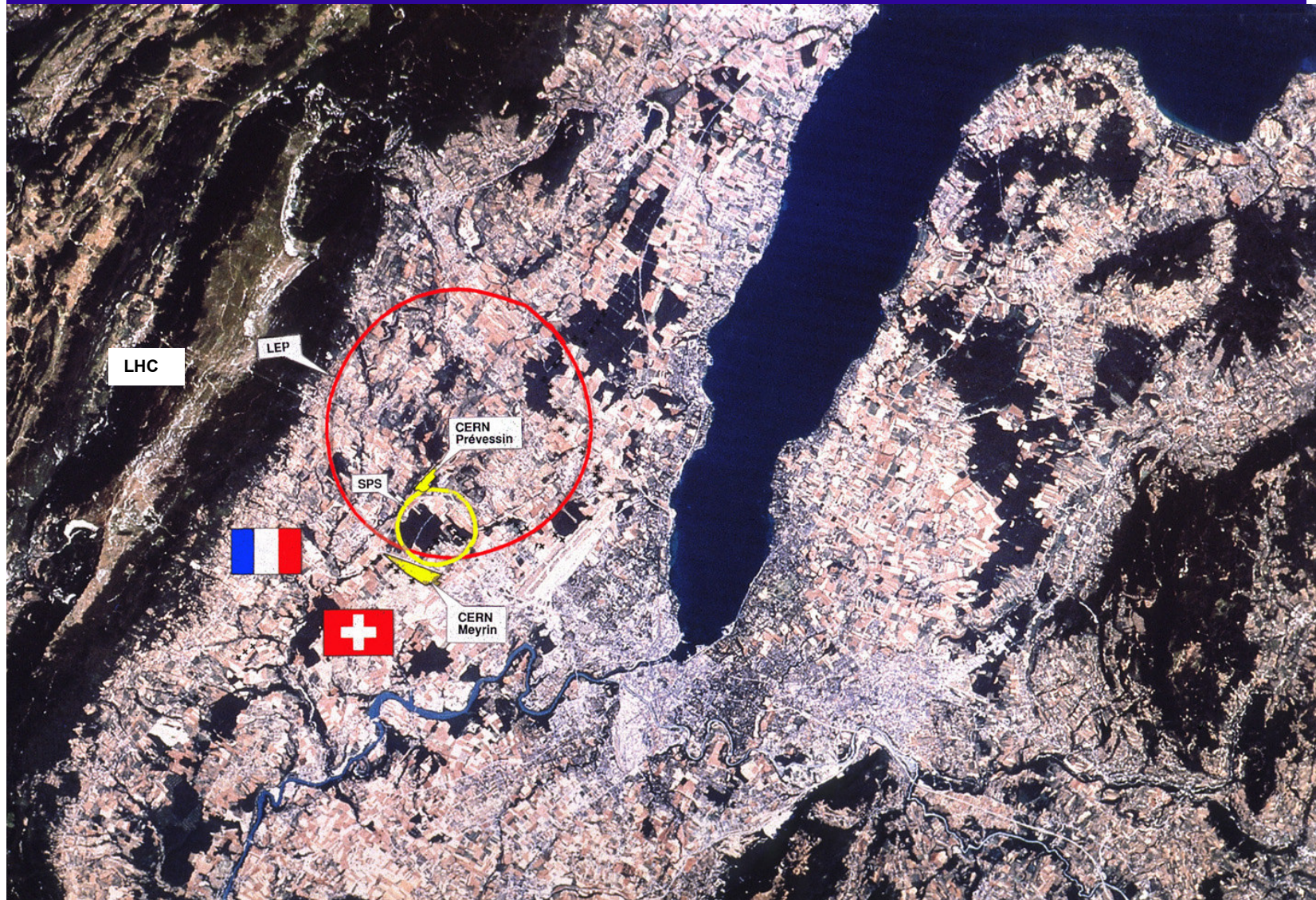


Dark Matter Experiments and Searches

R.J.Cashmore
Principal Brasenose College, Oxford
and
Dept of Physics, Oxford

Dark Matter at LHC

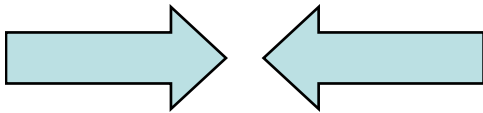
Satellite view of Geneva and CERN site



The **L**arge **H**adron **C**ollider in the LEP Tunnel

Proton- Proton Collider

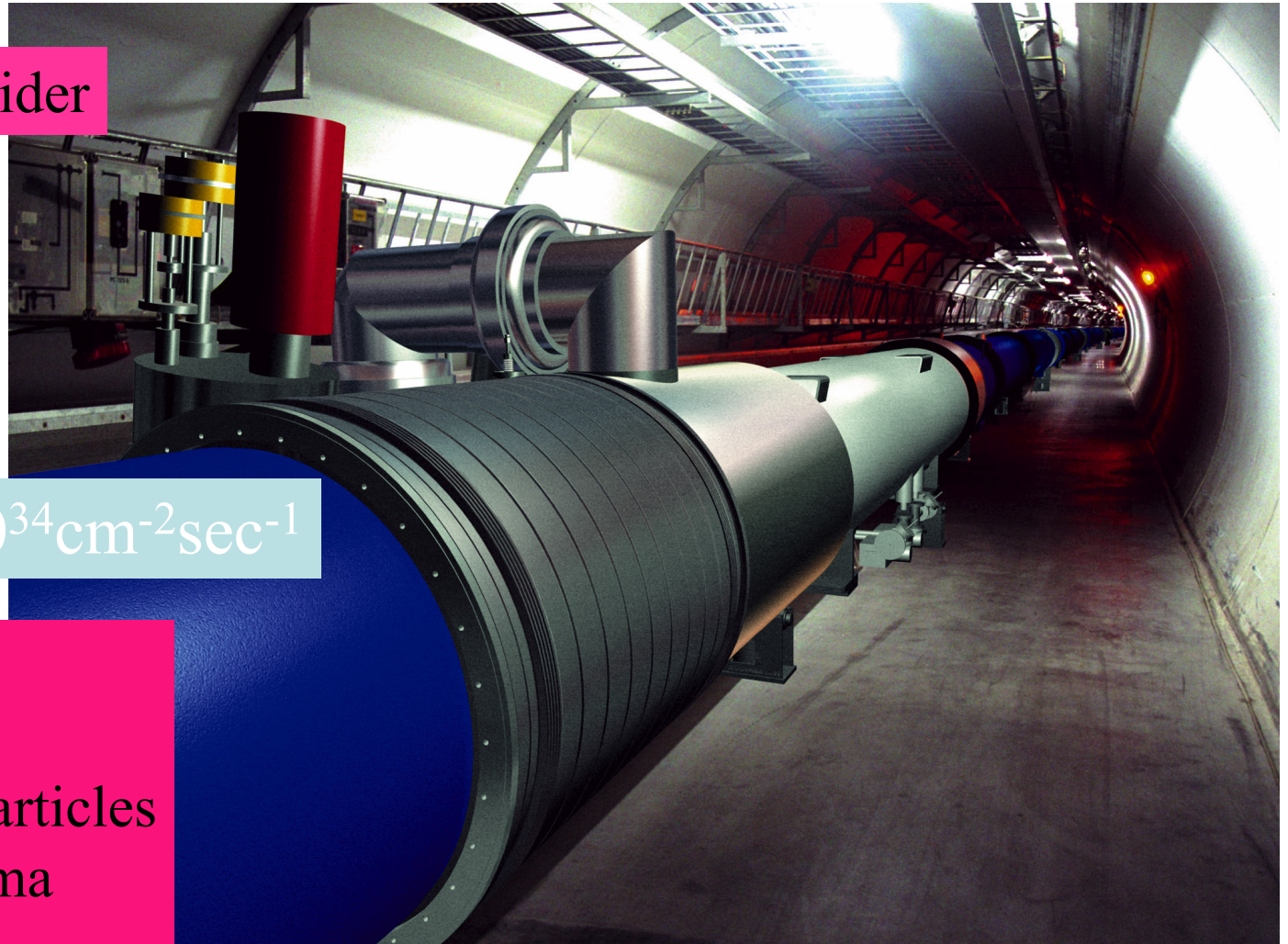
7 TeV + 7 TeV



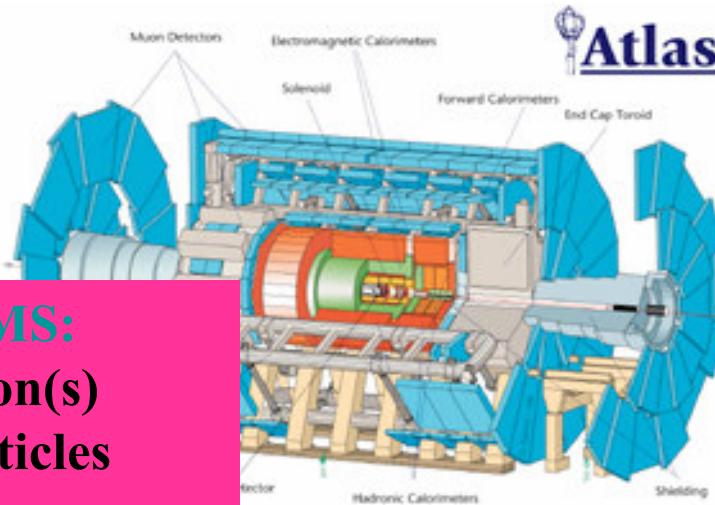
Luminosity = $10^{34} \text{cm}^{-2} \text{sec}^{-1}$

The Physics:

- Higgs boson (s)
- Supersymmetric Particles
- Quark-Gluon Plasma
- CP violation in B



LHC Experiments



ATLAS, CMS:

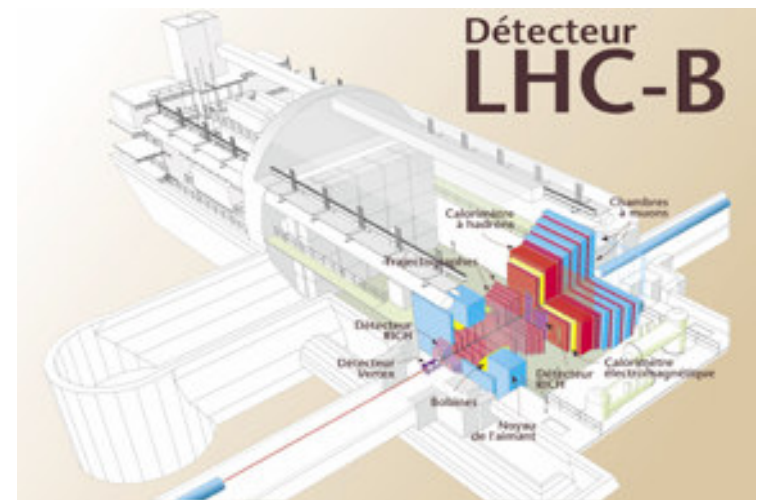
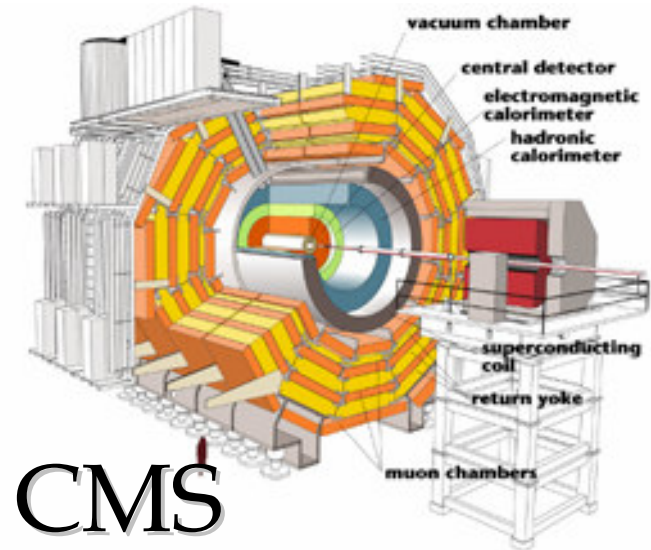
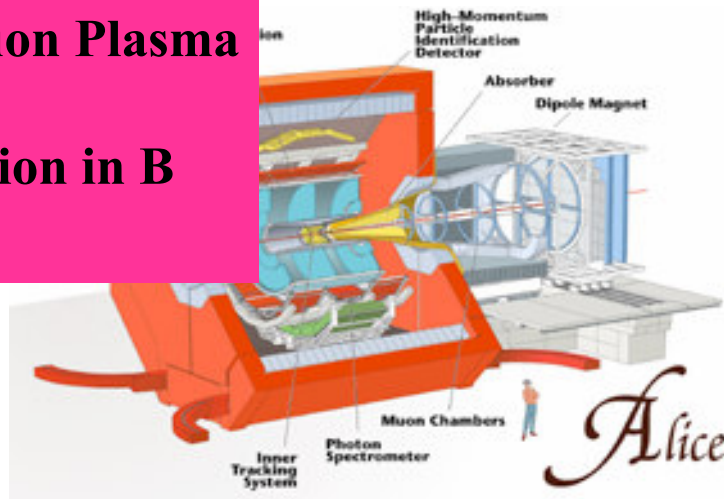
- Higgs boson(s)
- SUSY particles
- ...??

ALICE:

Quark Gluon Plasma

LHC-B:

- CP violation in B



Weakly Interacting Massive Particles

- The lightest of these new particles is an ideal candidate for dark matter
- Would have been created in the Big Bang
- A *natural* solution – makes it extremely attractive!

Now we need to detect these particles!

Particle type	Particle	Super partner
Fermion	Quark	Squark
	Neutrino	Sneutrino
	Electron	Selectron
	Muon	Smuon
	Tau	Stau
Boson	W	Wino
	Z	Zino
	Photon	Photino
	Gluon	Gluino
	Higgs	Higgsino

SUSY - sparticles

SUSYMSSM and mSUGRA parameter space
cross sections
decay BRs
SUSY event topology and mass reach
sparticle reconstruction
SUSY DM particles

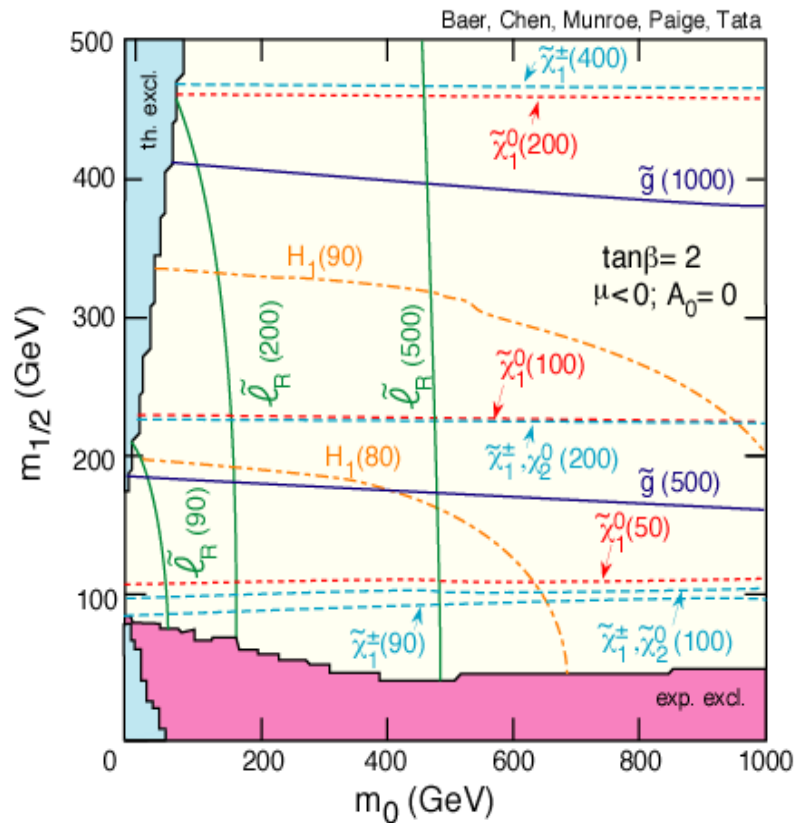
MSSM

- Two mass parameters: m_0 , $m_{1/2}$, one sign: μ
- Region in parameter space restricted by
 - non-observation of SUSY signal at LEP and Tevatron
 - Cosmological parameters reflecting cold dark mass properties
- Two Higgs doublets, physical scalar bosons:
 - Neutral: h , H , A (m_A and $\tan\beta$ can be used as independent parameters)
 - Charged: H^\pm
- SUSY partners of known particles:
 - spin 0: squarks, sleptons,
 - spin 1/2: gluinos, gauginos, Higgsino
- R-parity conservations: decay chains contain several neutral, invisible particles \rightarrow missing energy .

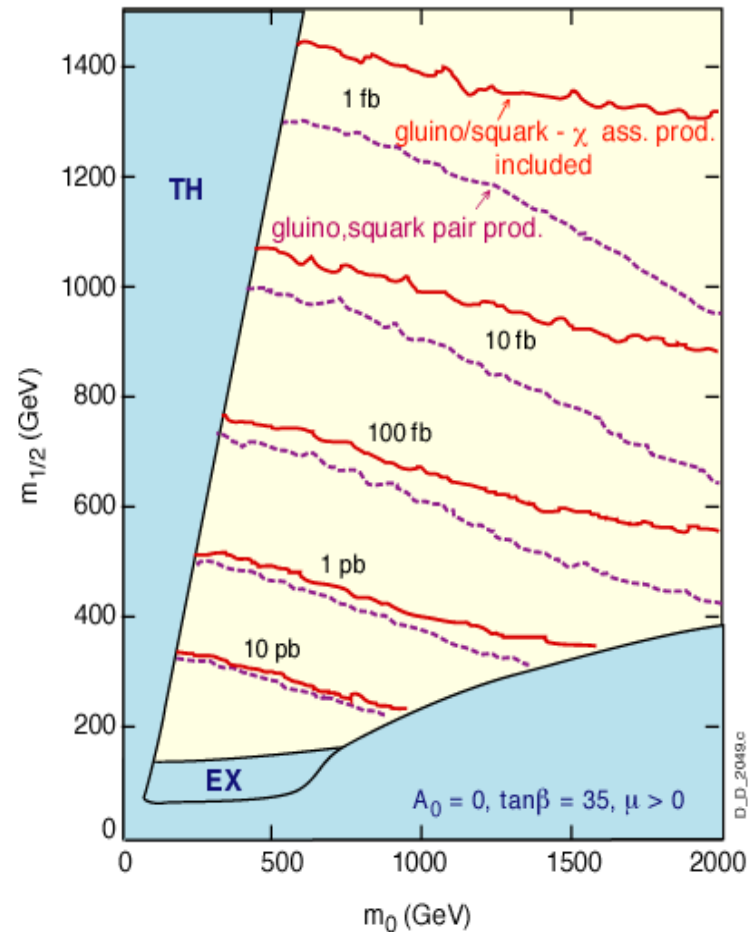
SUSY - parameter space, cross sections

Isomass contours in Minimal Supergravity, in m_0 vs $m_{1/2}$ plane

Parameters of theory: $m_0, m_{1/2}, A_0, \tan\beta, \text{sign}(\mu)$



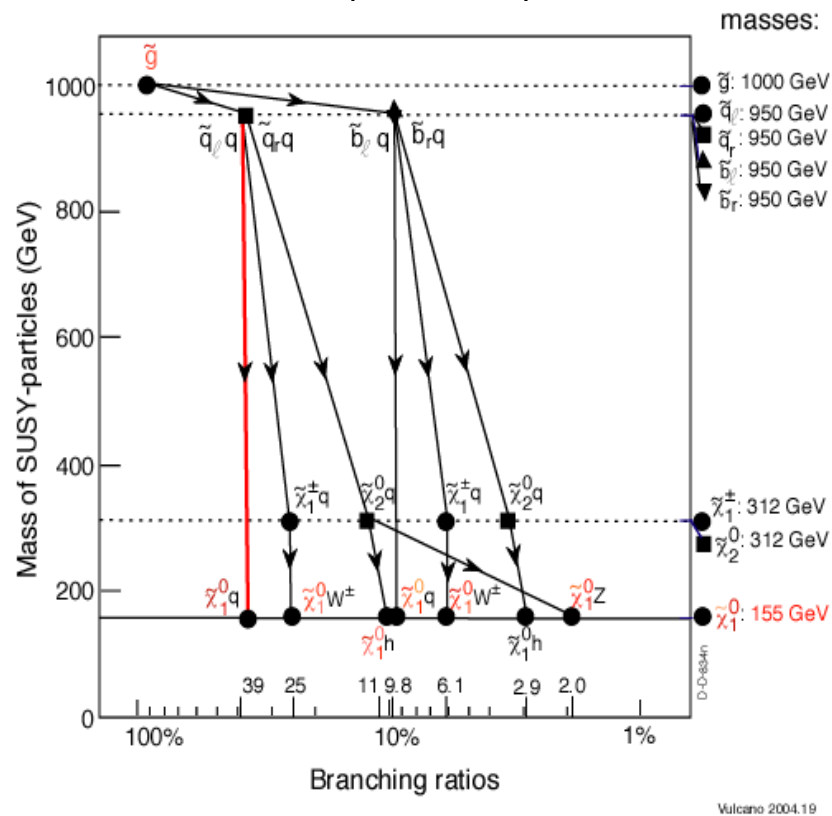
SUSY total cross sections (mSUGRA)



Glauino/squark decay chains

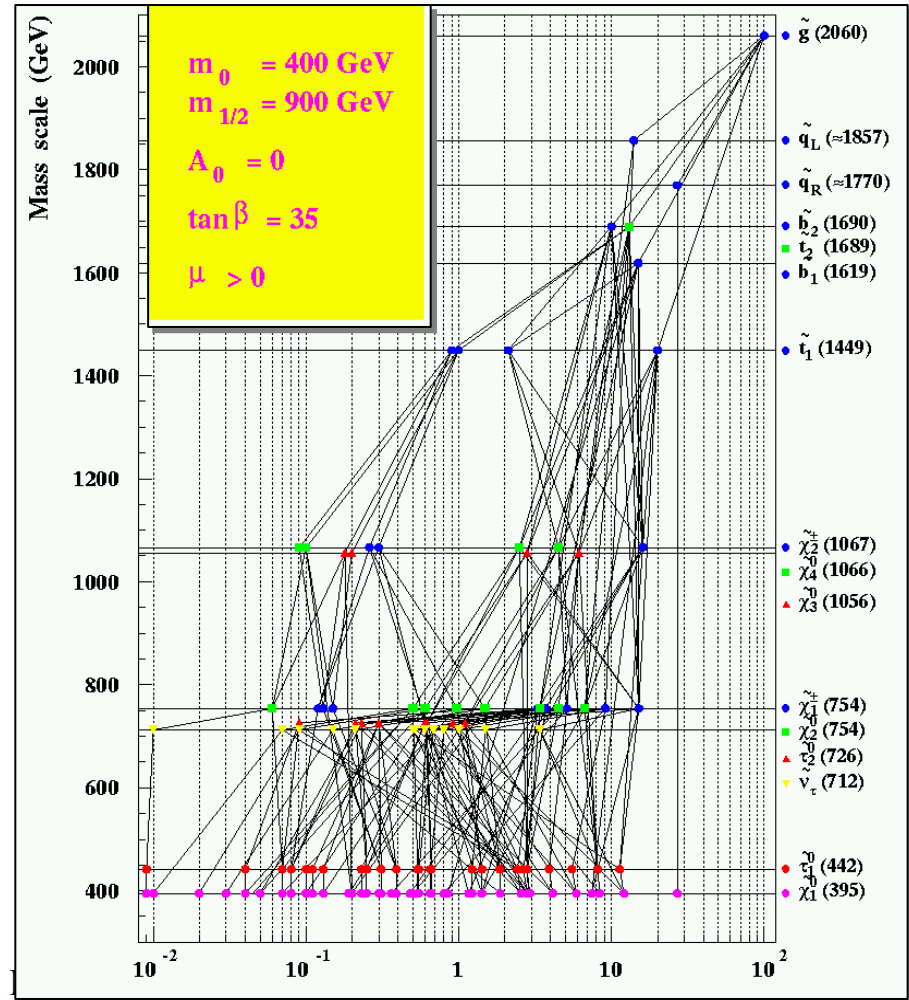
Typical decay chain for a massive gluino or squar
(mSUGRA)

Decay chain for a light gluino or squark
(MSSM)



R.Cashmore

Dark



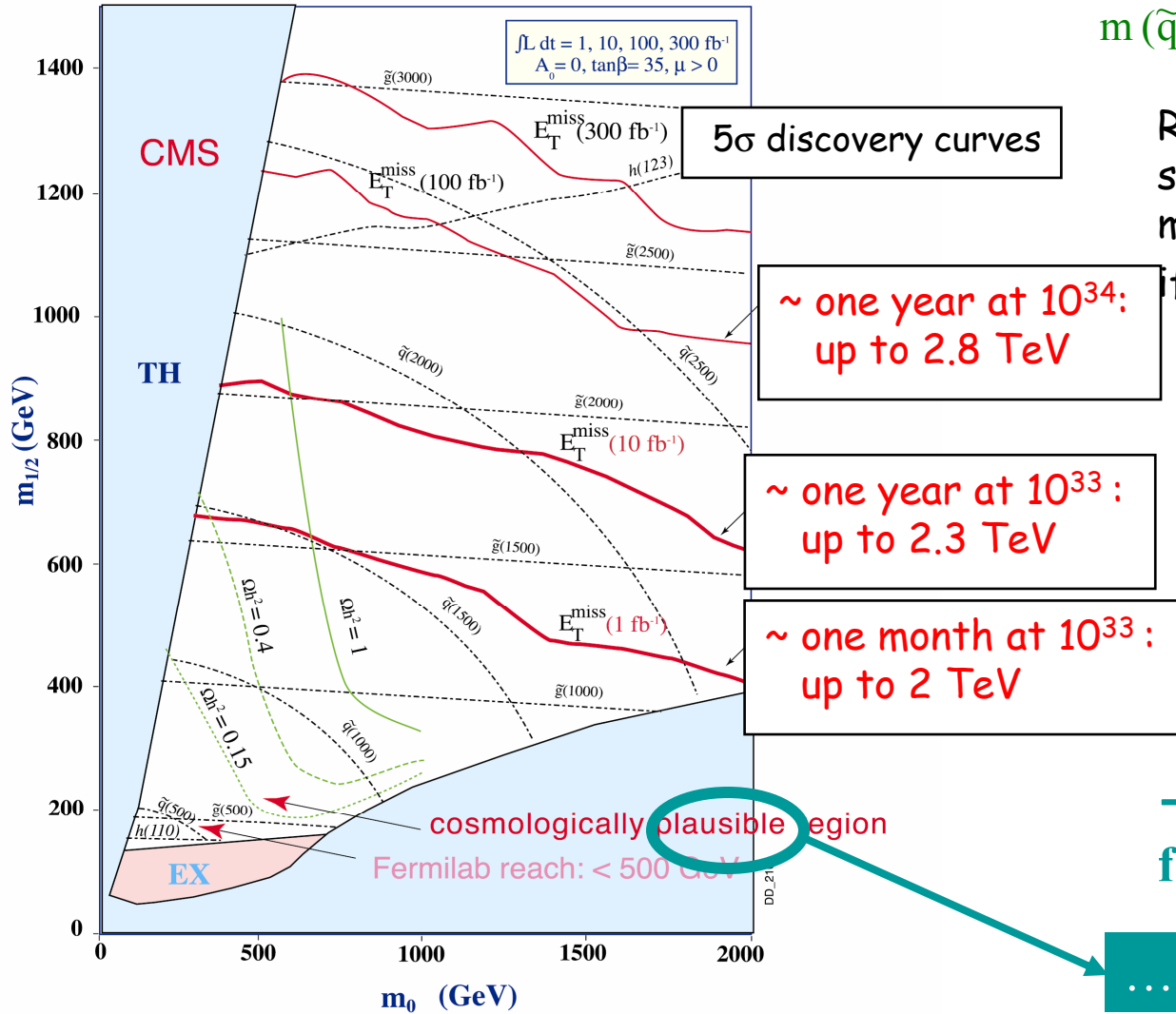
Vulcano 2004.6

SUPERSYMMETRY

Large $\tilde{q}\tilde{q}, \tilde{q}\tilde{g}, \tilde{g}\tilde{g}$

cross-section $\rightarrow \approx 100$ events/day at 10^{33} for:

$m(\tilde{q}, \tilde{g}) \sim 1$ TeV



Reach of **Multijet + E_T^{miss}** searches (most powerful and model-independent signature if R-parity conserved)

5 σ discovery curves

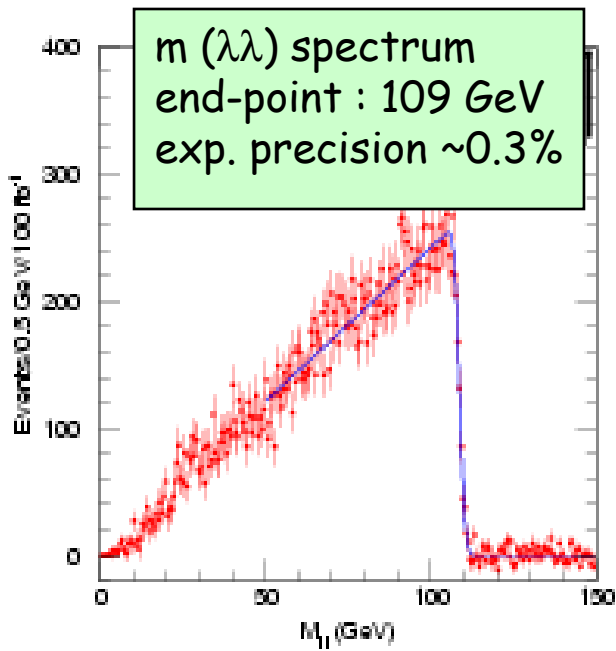
~ one year at 10^{34} : up to 2.8 TeV

~ one year at 10^{33} : up to 2.3 TeV

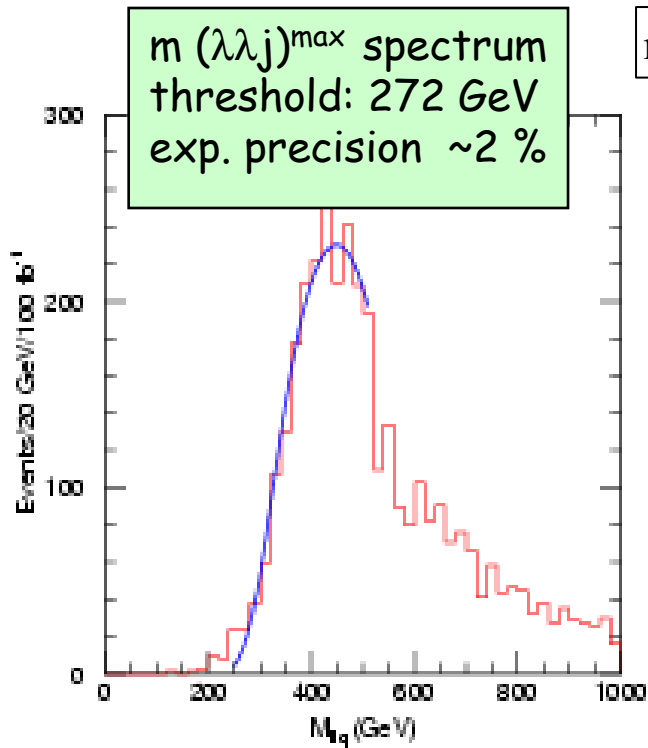
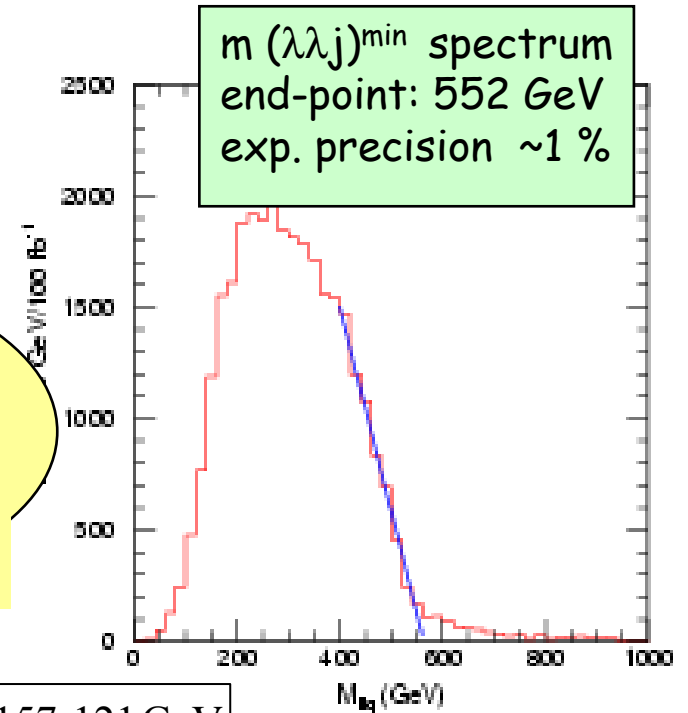
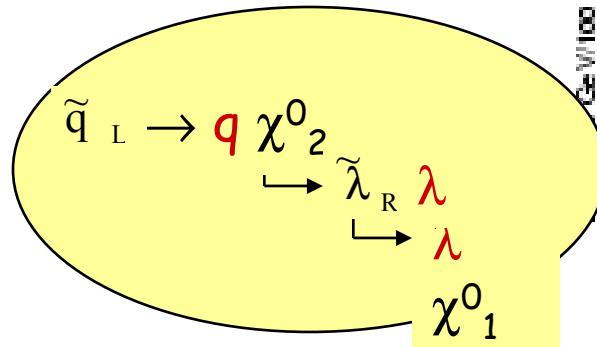
~ one month at 10^{33} : up to 2 TeV

\rightarrow SUSY could be found quickly

... “plausible” = typical



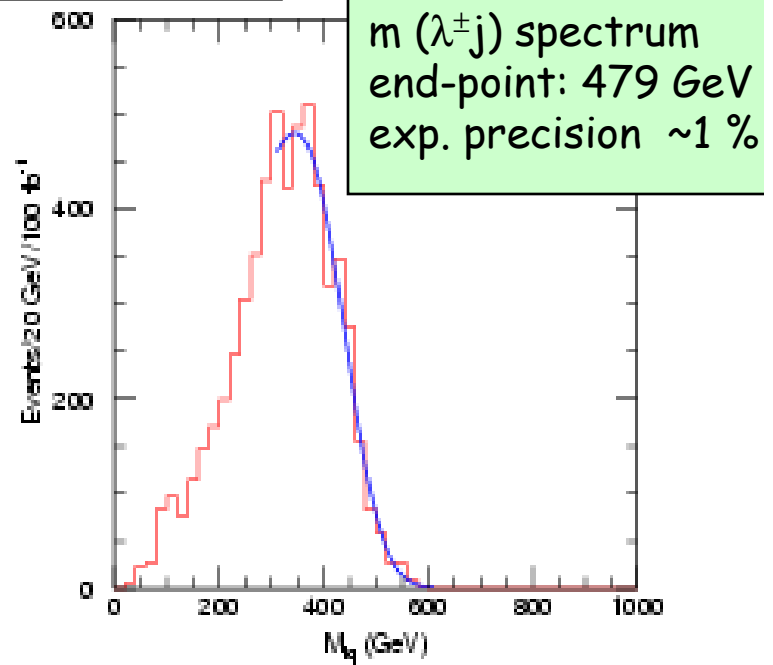
Example of
a typical chain:

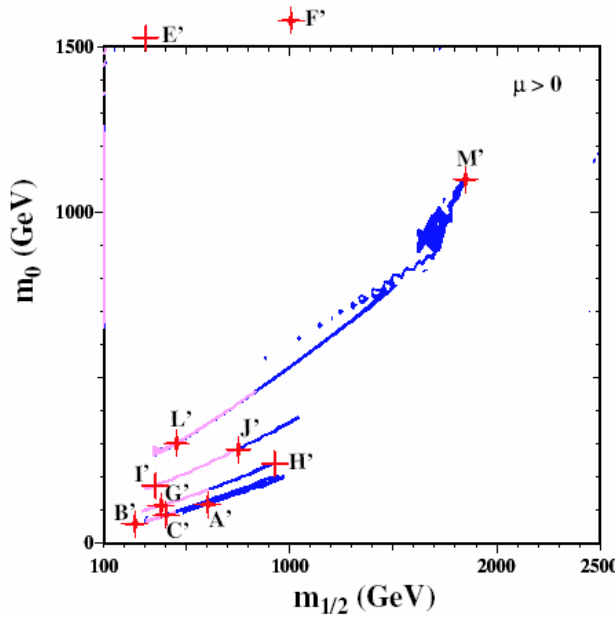


$$m(\tilde{q}_L, \chi^0_2, \tilde{\lambda}_R, \chi^0_1) = 690, 232, 157, 121 \text{ GeV}$$

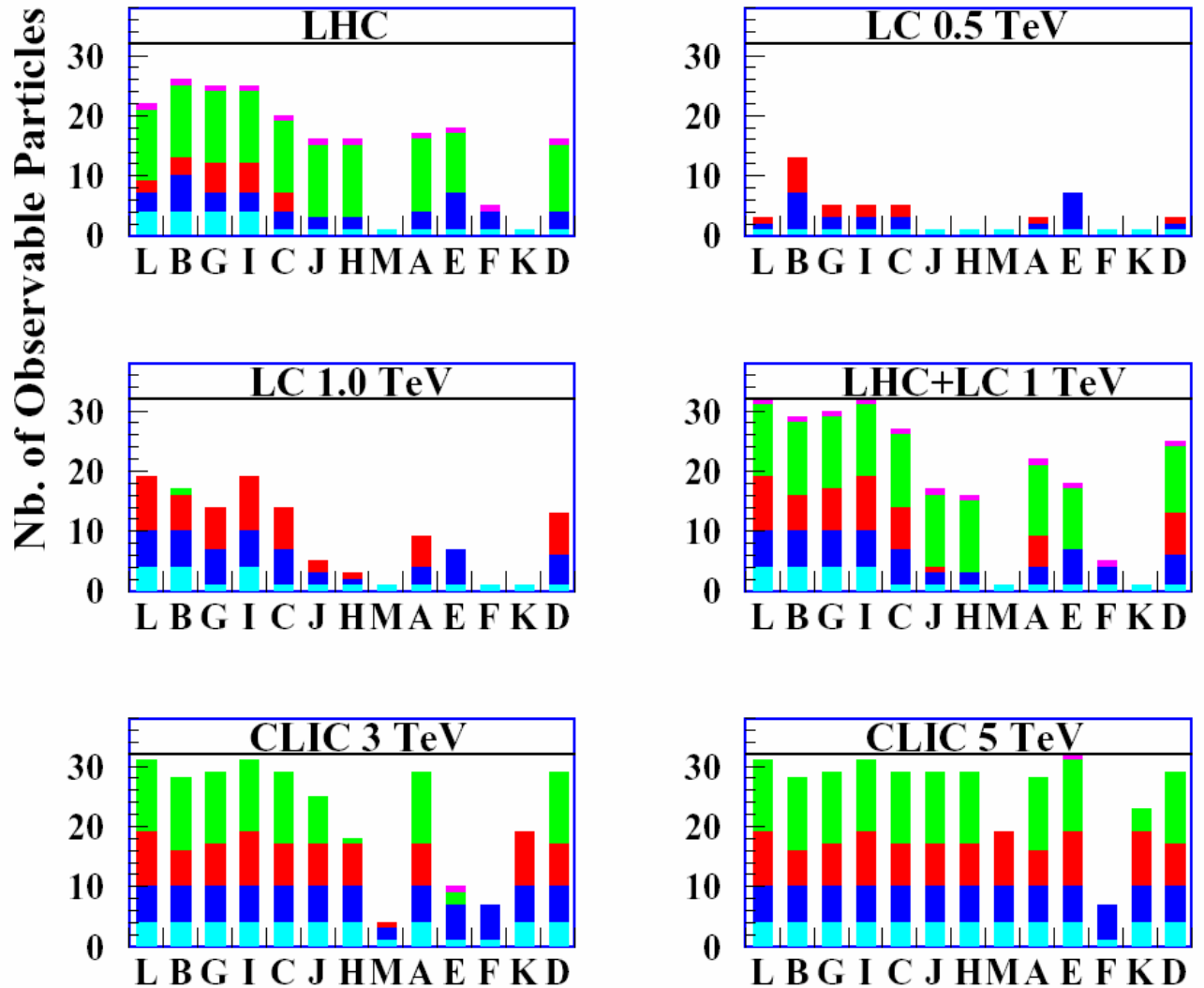
ATLAS
100 fb⁻¹
LHC Point 5

Dark Matter 3





█ gluino █ squarks █ sleptons █ χ █ H
Post-WMAP Benchmarks

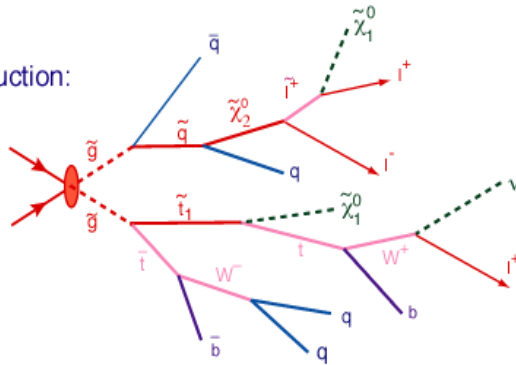


Benchmarking
 MSSMs restricted
 by Cosmological
 & Particle physics
 data
 M. Battaglia et al.

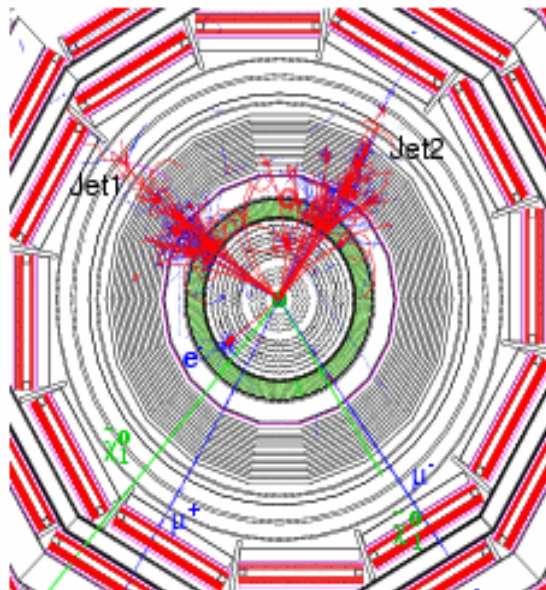
R. Cashmore

Squark and gluino searches, reach in various topologies

- Production:

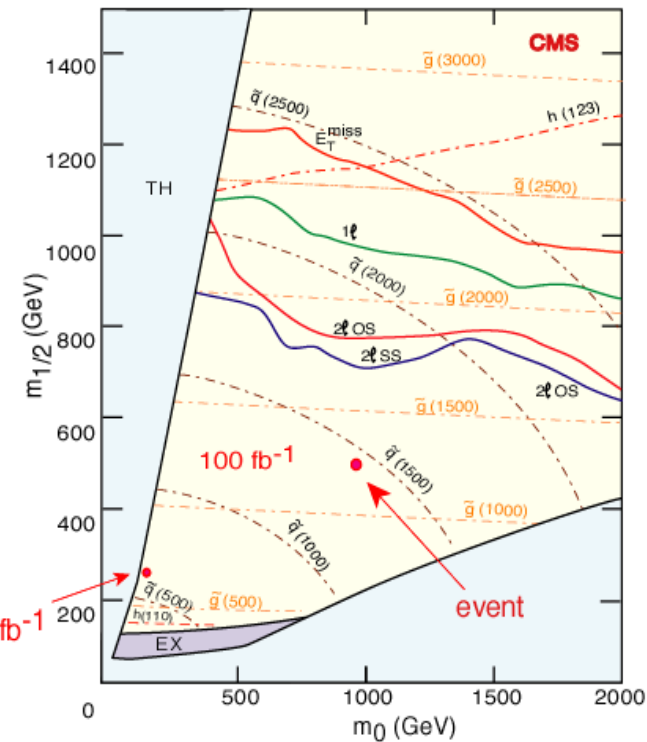


- More generally $\tilde{g}\tilde{g}$, $\tilde{g}\tilde{q}$, $\tilde{q}\tilde{q}$ production leads to: lepton(s) + E_T^{miss} + jets final states
- Backgrounds: W + jets, Z+jets, $t\bar{t}$, Wtb, WW, WZ..



investigate sparticle reconstruction for 10 fb^{-1} (but $\tan\beta = 10$)

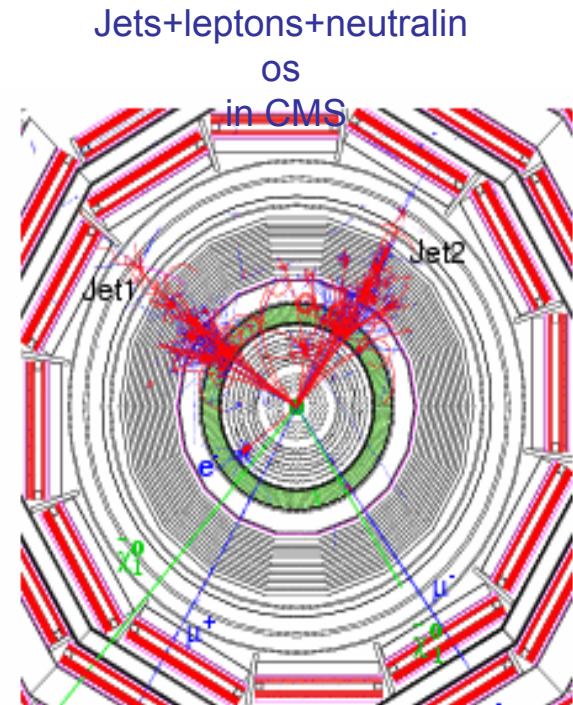
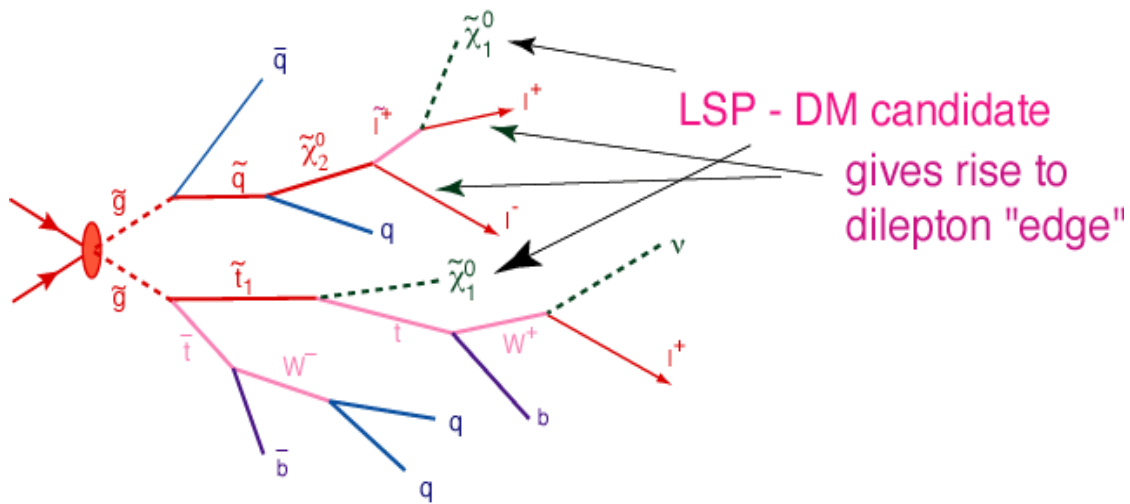
Explorable domain of m_0 , $m_{1/2}$ parameter space in \tilde{q}, \tilde{g} searches in n leptons + $E_T^{\text{miss}} + \geq 2$ jets
 m SUGRA, $A_0 = 0$, $\tan\beta = 35$, $\mu > 0$
 5σ contours; non-isolated muons



- Jets + E_T^{miss} has highest reach
- opposite-sign two-leptons are most useful for sparticle reconstruction

Sparticle reconstruction

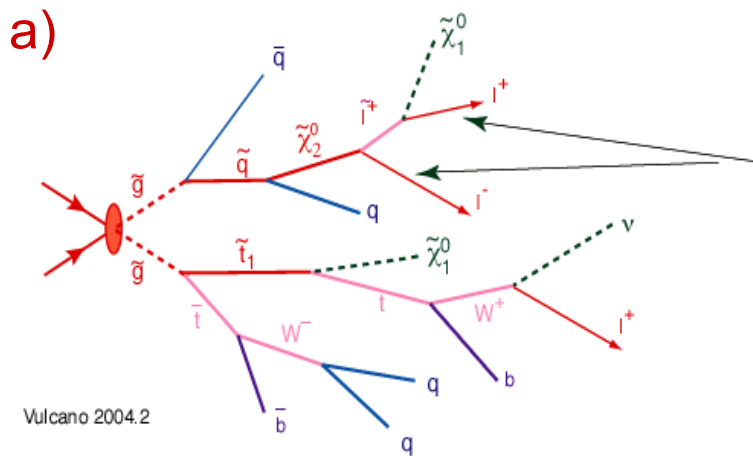
Gluino/squark production event topology allowing sparticle mass reconstruction



Such cascade decays allow to reconstruct sleptons, neutralinos, squarks, gluinos...in favorable cases with %level mass resolutions

Starting point of reconstruction is the "dilepton edge" due to the neutralino-2 decays to dileptons + neutralino-1;
neutralino-2 decay to neutralino-1 + h(\rightarrow bb) can be used too

Dilepton edges, sparticle reconstruction

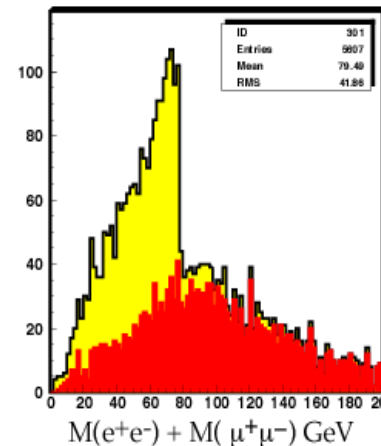


gives rise to dilepton "edge"

Starting point of reconstruction is the "dilepton edge" due to the neutralino-2 decays to dileptons + neutralino-1

b) m inclusive isolated dileptons + E_t^{miss} search

- $\tilde{q}, \tilde{g} \rightarrow \tilde{\chi}_2^0 + X$ with large BR
- $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$ has significant BR



Dilepton edge

$E_t^{\text{miss}} > 100$ GeV

SUSY point B

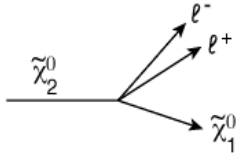
$m_{1/2} = 250$ GeV,

$m_0 = 100$ GeV,

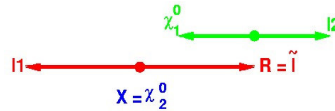
$\tan\beta = 10, \mu > 0$

c)

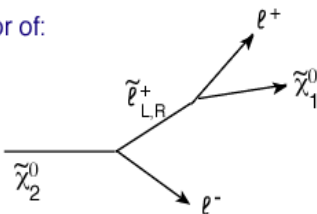
use kinematics of:



$$\Rightarrow M_{\ell\ell}^{\text{max}} = m(\tilde{\chi}_2^0) - m(\tilde{\chi}_1^0)$$



or of:



$$\Rightarrow M_{\ell\ell}^{\text{max}} = m(\tilde{\chi}_2^0) \sqrt{1 - \frac{m^2(\tilde{\ell})}{m^2(\tilde{\chi}_2^0)}} \sqrt{1 - \frac{m^2(\tilde{\chi}_1^0)}{m^2(\tilde{\ell})}}$$

d)

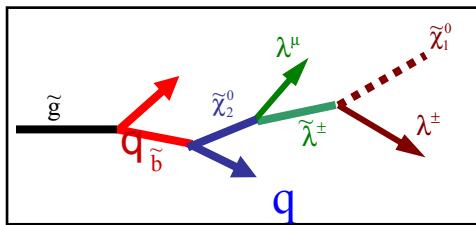
on the edge of $\tilde{\chi}_2^0 \rightarrow \ell^+ \ell^- \tilde{\chi}_1^0$:

$$\vec{P}_{\tilde{\chi}_2^0} = (1 + M_{\tilde{\chi}_1^0} / M_{\ell^+ \ell^-}) \vec{P}_{\ell^+ \ell^-} \text{ in lab. frame}$$

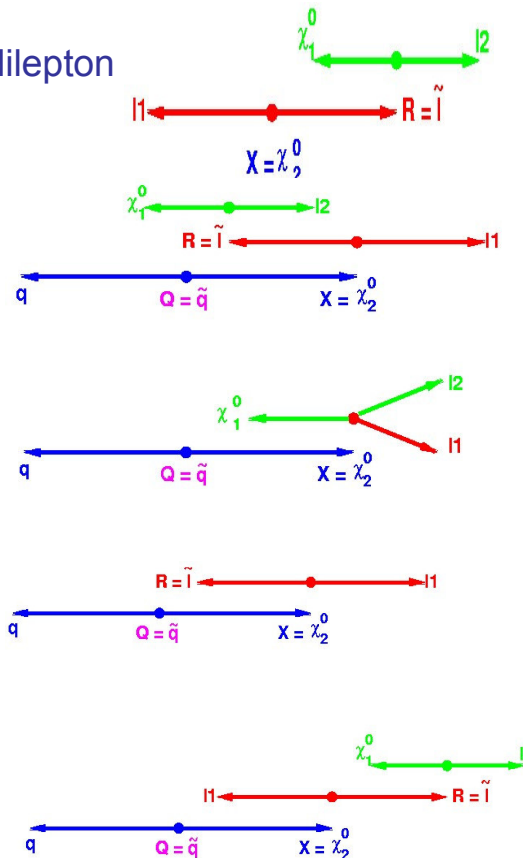
\Rightarrow reconstruct $\tilde{q}_L \rightarrow \tilde{\chi}_2^0 q$ decay combining $\tilde{\chi}_2^0$ with q-jet and get $m(\tilde{q})$, assuming $m(\tilde{\chi}_1^0) = 0.5 m(\tilde{\chi}_2^0)$ (for ex.)

Sparticle reconstruction, end-points

In a long cascade of decays a number of particular kinematical configurations generate end-points depending on sparticle masses, fitting these end points allows to constrain or determine these masses



for ex.:
largest dilepton
mass:



Dark Matter 3

End-points depend on mass differences, thus strong correlations between masses;
A fit of mSUGRA predictions to the set of measured end-points can be performed, this yields model parameters and sparticle masses, neutralino mass and relic density, then check for consistency with terrestrial, CMB experiments.....

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$$M_{\lambda\lambda}^{\max} = \frac{\sqrt{(m_{\chi_2^0}^2 - m_{\tilde{\lambda}}^2)(m_{\tilde{\lambda}}^2 - m_{\chi_1^0}^2)}}{m_{\tilde{\lambda}}}$$

$$M_{\lambda\lambda q}^{\max} = \frac{\sqrt{(m_{\chi_2^0}^2 - m_{\chi_1^0}^2)(m_{\tilde{q}}^2 - m_{\chi_2^0}^2)}}{m_{\chi_2^0}}$$

$$M_{\lambda\lambda q}^{\max} = m_{\tilde{q}} - m_{\chi_1^0} \text{ if } m_{\chi_2^0}^2 > m_{\tilde{q}} m_{\chi_1^0}$$

$$M_{\lambda_1 q}^{\max} = \frac{\sqrt{(m_{\chi_2^0}^2 - m_{\tilde{\lambda}}^2)(m_{\tilde{q}}^2 - m_{\chi_2^0}^2)}}{m_{\chi_2^0}}$$

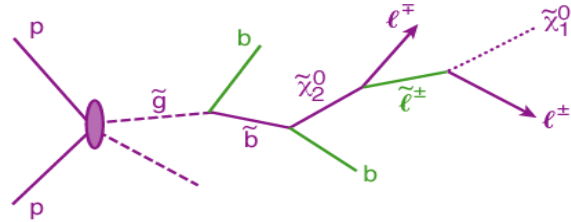
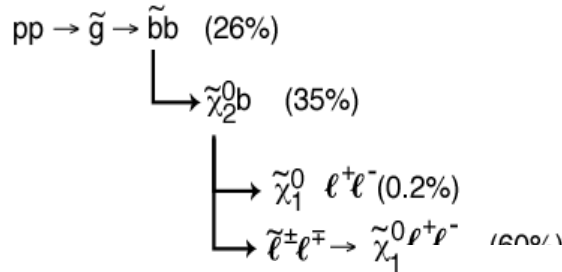
$$M_{\lambda_2 q}^{\max} = \frac{\sqrt{(m_{\tilde{\lambda}}^2 - m_{\chi_1^0}^2)(m_{\tilde{q}}^2 - m_{\chi_2^0}^2)}}{m_{\tilde{\lambda}}}$$

Sparticle reconstruction, full chain

Glino and sbottom reconstruction

Post-LEP SUSY benchmark point B
 $(m_{1/2} = 250 \text{ GeV}, m_0 = 100 \text{ GeV}, \tan\beta = 10, \mu > 0)$

Decay chain and event topology:

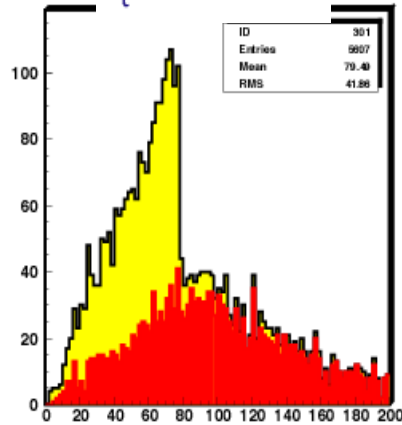


Event final state:

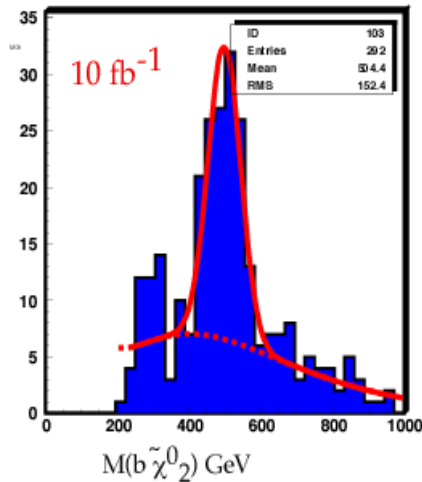
- ≥ 2 high p_T isolated leptons OS
- ≥ 2 high p_T b jets
- missing E_T

Dilepton edge

$E_T^{\text{miss}} > 100 \text{ GeV}$



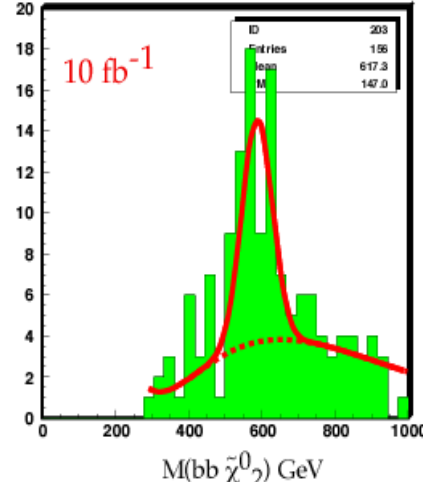
Sbottom reconstruction



fit:
 $M(b \tilde{\chi}_2^0) \text{ GeV} = 499 \pm 7 \text{ GeV}$
 $\sigma = 48 \text{ GeV}$

generated: $M(\tilde{b}_L) = 496 \text{ GeV}$
 $M(\tilde{b}_R) = 524 \text{ GeV}$

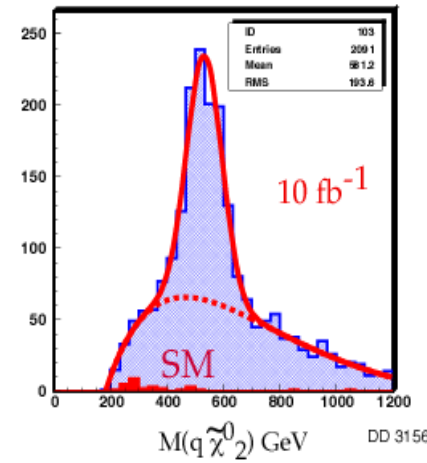
Glino reconstruction



fit:
 $M(bb \tilde{\chi}_2^0) \text{ GeV} = 587 \pm 11 \text{ GeV}$
 $\sigma = 41 \text{ GeV}$

generated mass: $M(\tilde{g}) = 595 \text{ GeV}$

Squark reconstruction anti-b-tag: $\sigma < 2$, b-veto



fit:
 $M(q \tilde{\chi}_2^0) = 532 \pm 4 \text{ GeV}$
 $\sigma = 62 \text{ GeV}$

$M(\tilde{d}_L) = M(\tilde{s}_L) = 543 \text{ GeV}$
 $M(\tilde{u}_L) = M(\tilde{c}_L) = 537 \text{ GeV}$

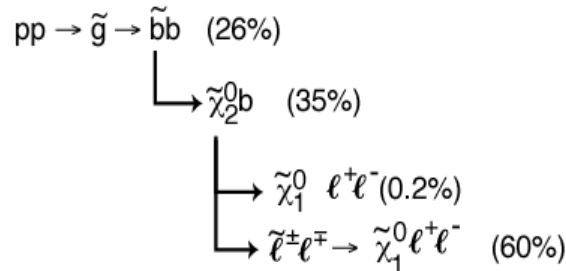
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Sparticle reconstruction, full chain

Glino and sbottom reconstruction

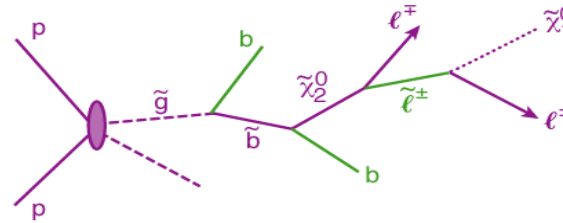
Post-LEP SUSY benchmark point B
 $(m_{1/2} = 250 \text{ GeV}, m_0 = 100 \text{ GeV}, \tan\beta = 10, \mu > 0)$

Decay chain and event topology:



Dilepton edge

$E_t^{\text{miss}} > 100 \text{ GeV}$



Event final state:

- ≥ 2 high p_t isolated leptons OS
- ≥ 2 high p_t b jets
- missing E_t

fit:

$$M(b \tilde{\chi}_2^0) \text{ GeV} = 499 \pm 7 \text{ GeV}$$

$$\sigma = 48 \text{ GeV}$$

generated:

$$M(\tilde{b}_L) = 496 \text{ GeV}$$

$$M(\tilde{b}_R) = 524 \text{ GeV}$$

fit:

$$M(bb \tilde{\chi}_2^0) \text{ GeV} = 587 \pm 11 \text{ GeV}$$

$$\sigma = 41 \text{ GeV}$$

generated mass: $M(\tilde{g}) = 595 \text{ GeV}$

fit:

$$M(q \tilde{\chi}_2^0) = 532 \pm 4 \text{ GeV}$$

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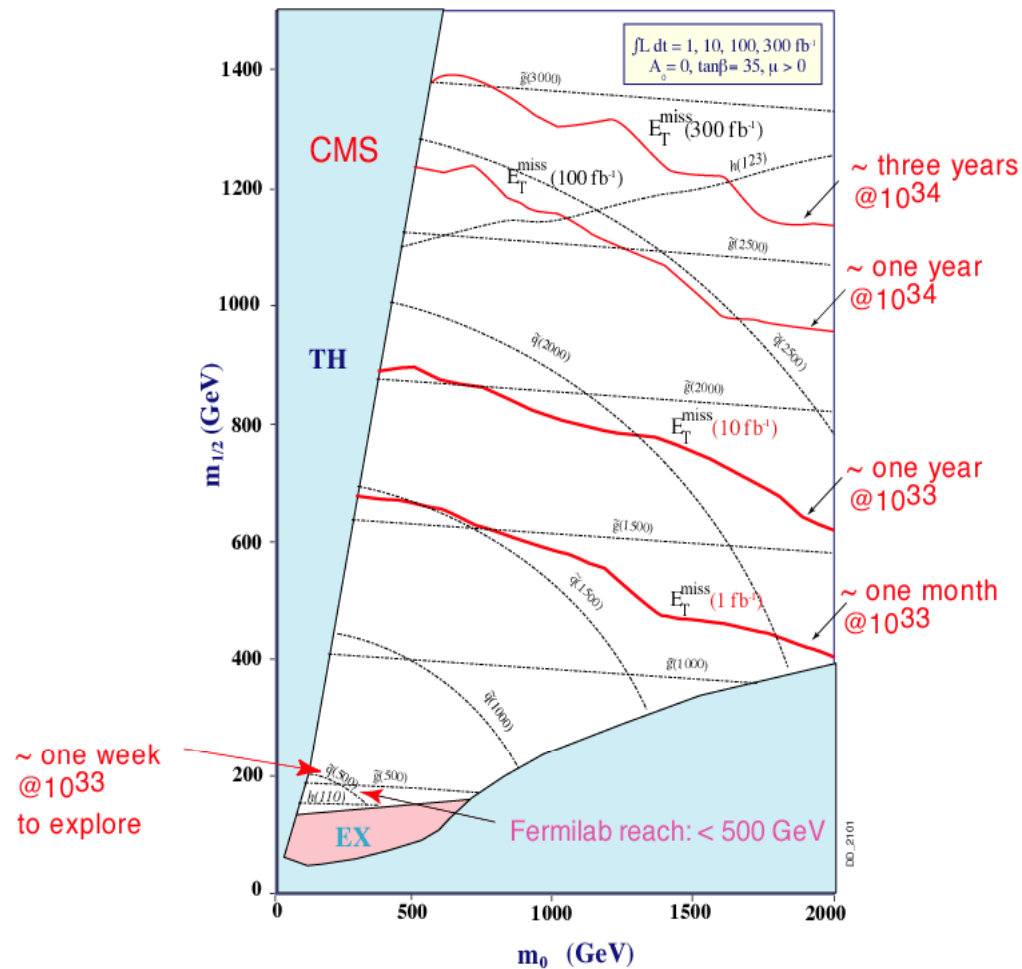
$$M(\tilde{d}_L) = M(\tilde{s}_L) = 543 \text{ GeV}$$

$$M(\tilde{u}_L) = M(\tilde{c}_L) = 537 \text{ GeV}$$

Squark/gluino mass reach vs luminosity

Reach vs luminosity

CMS \tilde{q}, \tilde{g} mass reach in $E_T^{\text{miss}} + \text{jets}$ inclusive channel
for various integrated luminosities



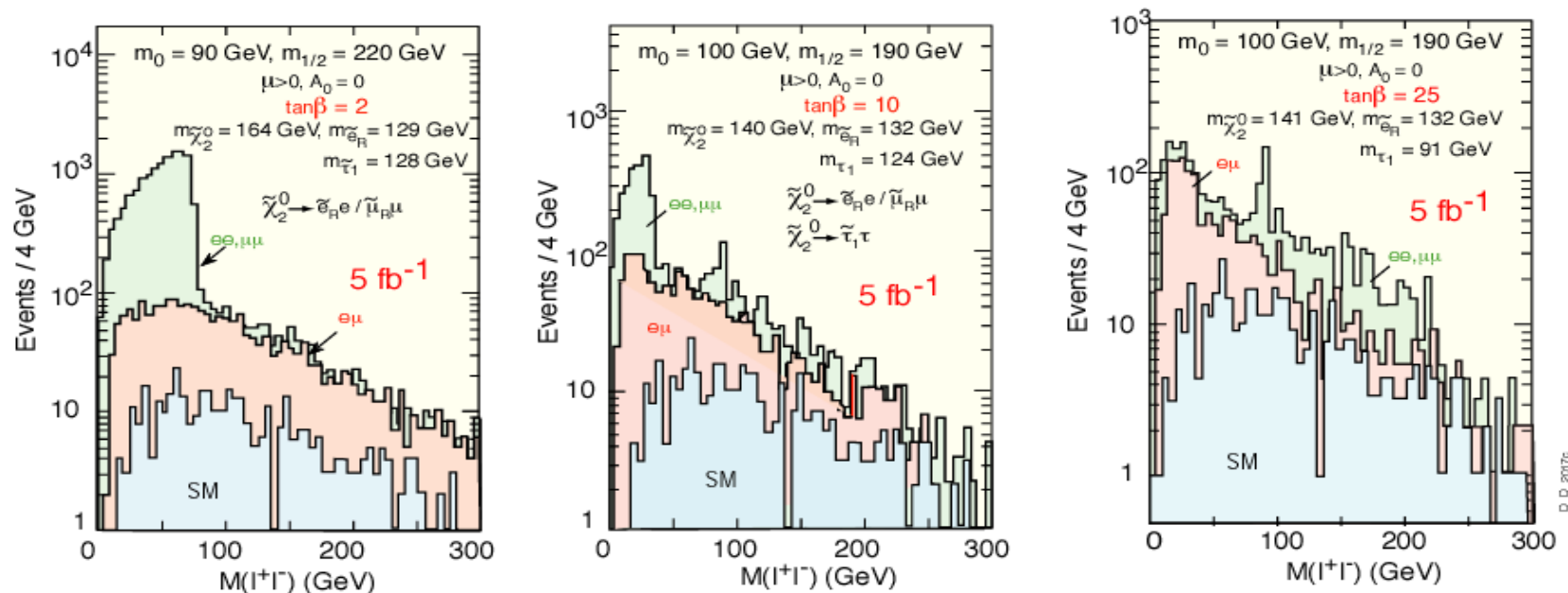
Dilepton edges at larger $\tan\beta$

Evolution of inclusive dilepton spectra with increasing $\tan\beta$

($m_0 \sim 100$ GeV, $m_{1/2} \sim 200$ GeV)

5 fb^{-1}

inclusive isolated opposite-sign dilepton spectra ($ee, \mu\mu, e\mu$)



→ SUSY may reveal itself early through peculiarities in dilepton spectra, for example such "edges" or "collapsed edges"

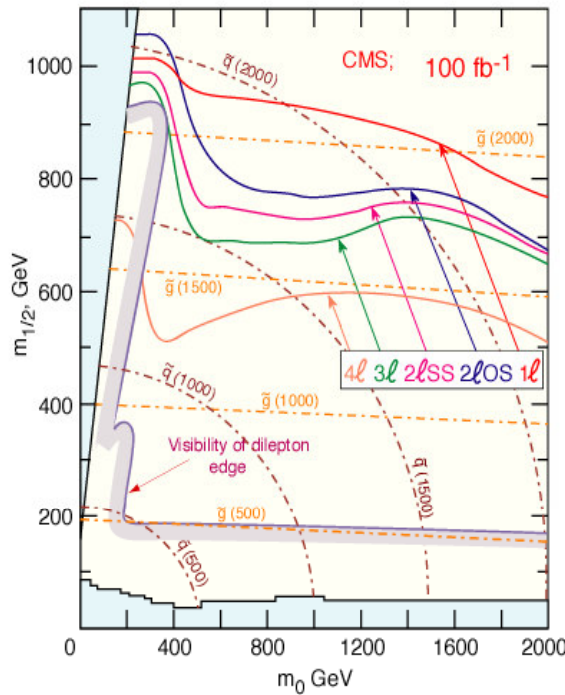
Observability of dilepton structures

Observability of " l^+l^- edge" or threshold structure in inclusive isolated two - leptons + E_t^{miss} final states

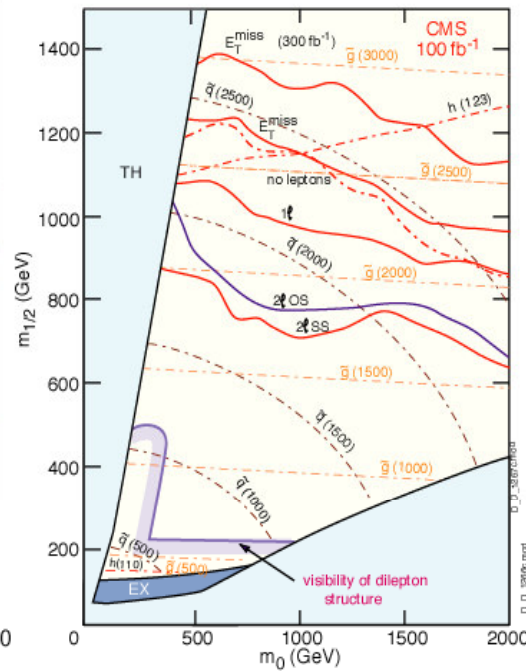
m SUGRA; $\tan \beta = 2$, $A_0 = 0$, $\mu < 0$
all leptons isolated

m SUGRA, $A_0 = 0$, $\tan \beta = 35$, $\mu > 0$
5 σ contours ; non - isolated muons

low $\tan \beta$ regime



high $\tan \beta$ regime



Structures tend to be less evident with increasing $\tan \beta$ and are limited to smaller regions of parameter space

Sparticle reconstruction, use several edges

Post-LEP SUSY benchmark point B, ATLAS/2004-007

A fit of mSUGRA predictions to the set of measured end-points can be performed

Edge	Nominal Value	Fit Value	Syst. Error Energy Scale	Statistical Error
$m(ll)^{\text{edge}}$	77.077	77.024	0.08	0.05
$m(ql)^{\text{edge}}$	431.1	431.3	4.3	2.4
$m(ql)_{\text{min}}^{\text{edge}}$	302.1	300.8	3.0	1.5
$m(ql)_{\text{max}}^{\text{edge}}$	380.3	379.4	3.8	1.8
$m(ql)^{\text{thres}}$	203.0	204.6	2.0	2.8
$m(bll)^{\text{thres}}$	183.1	181.1	1.8	6.3

+ eight other measurements

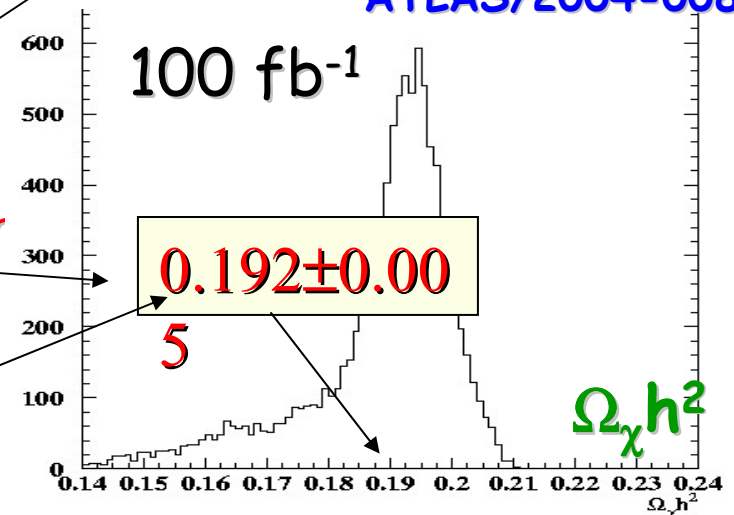
Parameter	Expected precision 30 fb ⁻¹	300 fb ⁻¹
m_0	± 3.2%	± 1.4%
$m_{1/2}$	± 0.9%	± 0.6%
$\tan(\beta)$	± 0.5%	± 0.5%

ATLAS/2004-008

Fit the best mSUGRA point by minimizing the overall χ^2

Deduce the LSP mass (with ~10% uncertainty) and relic density at the point

Determine $\Omega_\chi h^2$ and check for consistency with WMAP, terrestrial, astroparticle etc

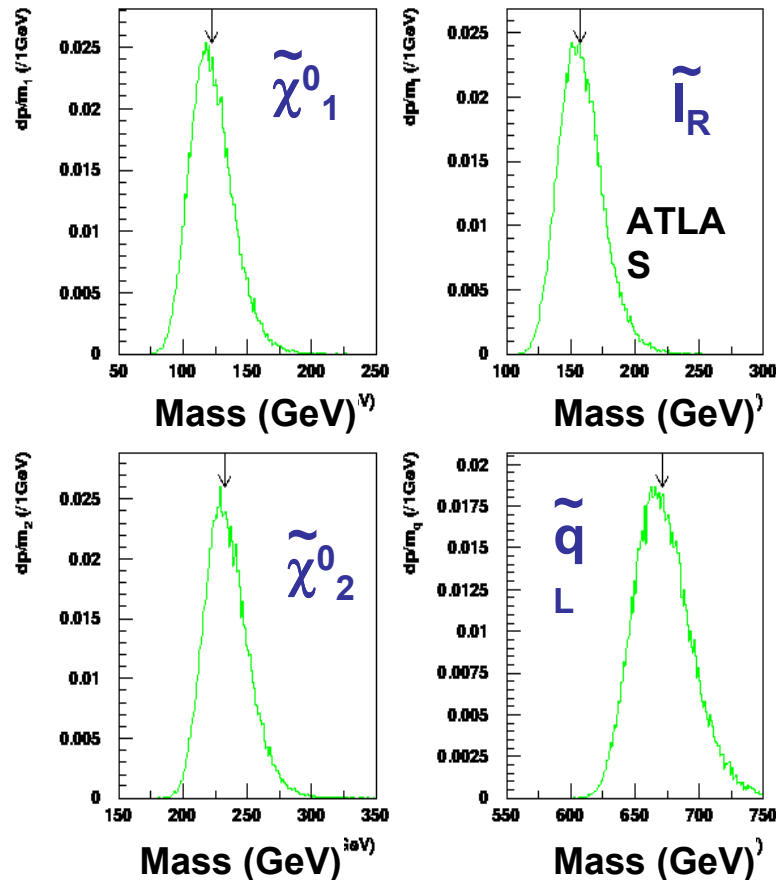


Model-independent predictions

Alternative approach to CMSSM fit to edge positions

numerical solution of simultaneous edge position equations

but interpretation of chain model dependent



ATLAS study results:

Sparticle	Expected precision (100 fb ⁻¹)
\tilde{q}_L	$\pm 3\%$
$\tilde{\chi}^0_2$	$\pm 6\%$
\tilde{l}_R	$\pm 9\%$
$\tilde{\chi}^0_1$	$\pm 12\%$

Use approximations together with other measurements to obtain 'model-independent' estimates of $\Omega_\chi h^2$, $\sigma_{\chi p}$, ϕ_{sun} etc.

Dark Matter, SUSY, LHC (I)

Dark matter:

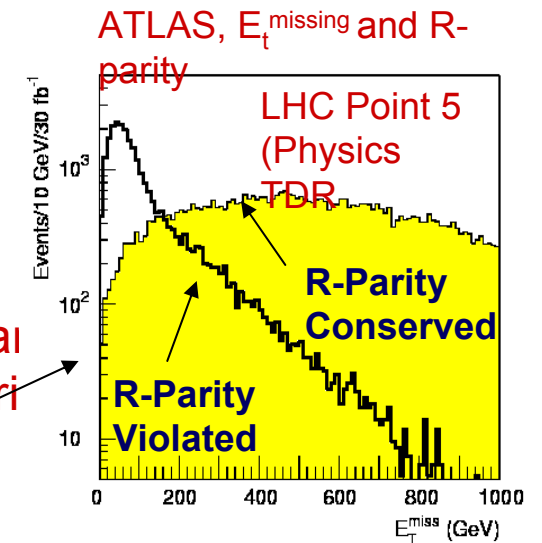
baryonic (machos, but represent no more than 20% of galactic halo DM),
 elementary particles, particle physics candidates:

axion,

neutralino (CDM), χ^0_1

neutrino (HDM), ν_τ , but too light, not more than
 few% of DM content

the neutralino-1 (χ^0_1) is a viable DM candidate only so far
 as the Lightest Supersymmetric Particle (LSP) and R-parity
 is conserved i.e. is a stable particle



DM detection:

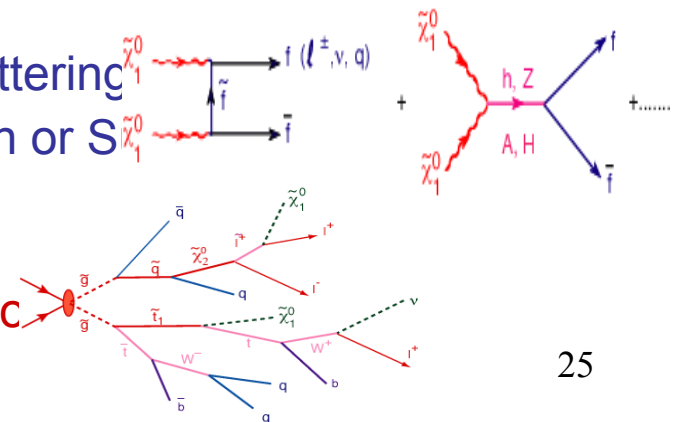
direct (ex.: relic neutralino (WIMP)-nucleon scattering)

indirect (ex. relic neutralino annihilation in Earth or S
 giving, for example, two neutrinos:

neutralino production, as at LHC (or the Tevatron)

R.Cashmore

Dark Matter 3



25

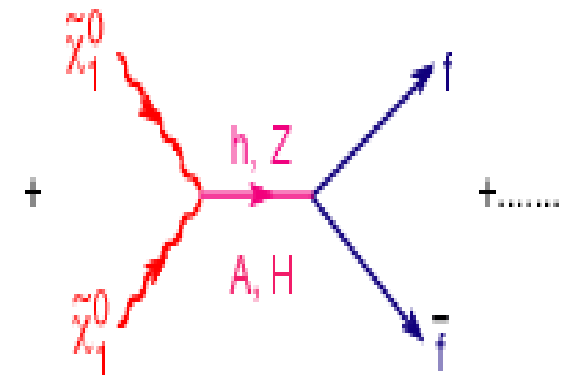
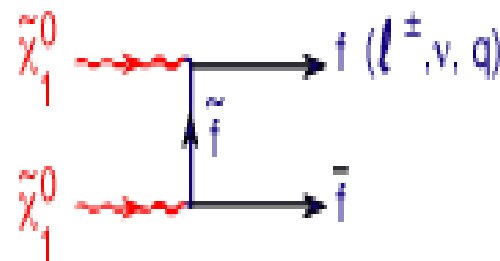
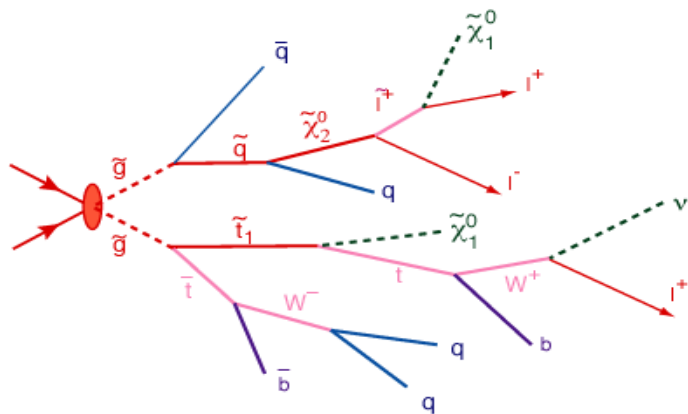
Dark Matter, SUSY, LHC (I)

DM detection:

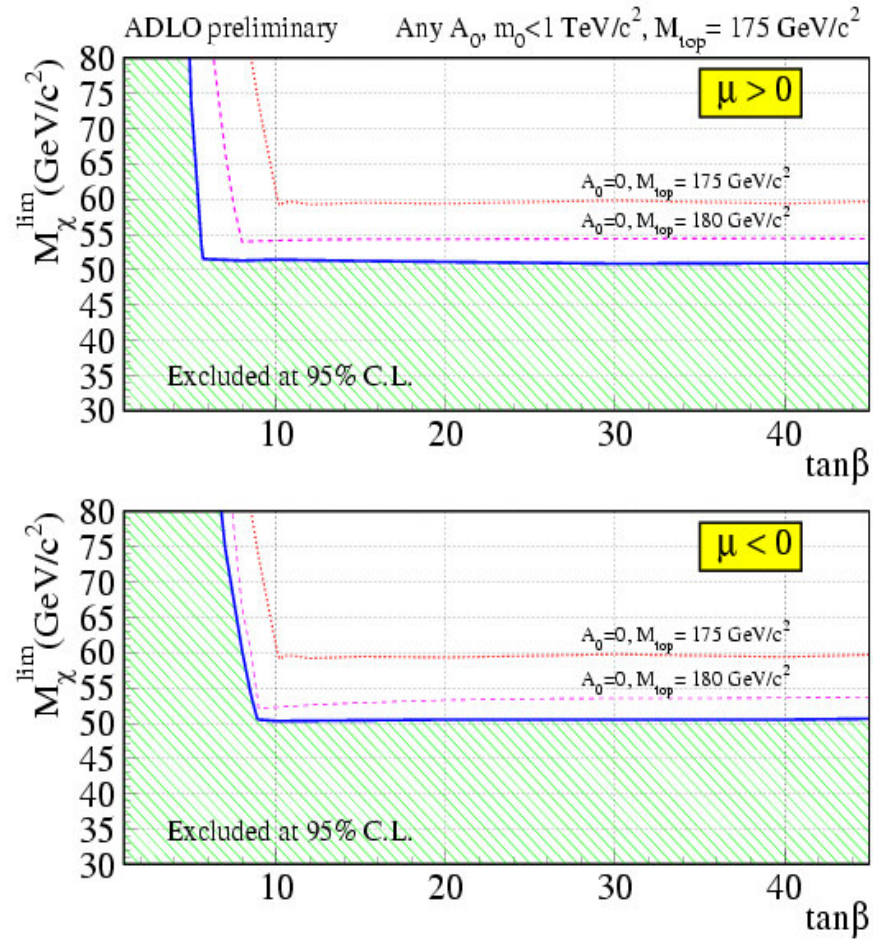
direct (ex.: relic neutralino (WIMP)-nucleon scattering)

indirect (ex. relic neutralino annihilation in Earth or Sun giving, for example, two neutrinos:

neutralino production, as at LHC (or the Tevatron):



LEP neutralino limits



Dark Matter, SUSY, LHC (II)

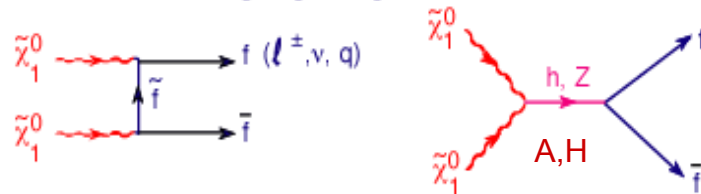
Two questions:

- do present (WMAP) constraints on DM assure SUSY detection at the LHC, and what limits do they imply on sparticle masses?
- can we measure the neutralino-1 mass?

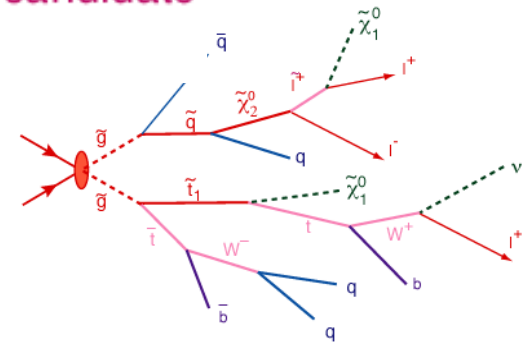
Relic densities in LHC mSUGRA parameter space (pre-WMAP)

The lightest neutralino $\tilde{\chi}_1^0$ is an excellent cold dark matter candidate (provided it is the LSP and R-parity is conserved)

Typical $\tilde{\chi}_1^0$ annihilations during Big Bang cooldown:

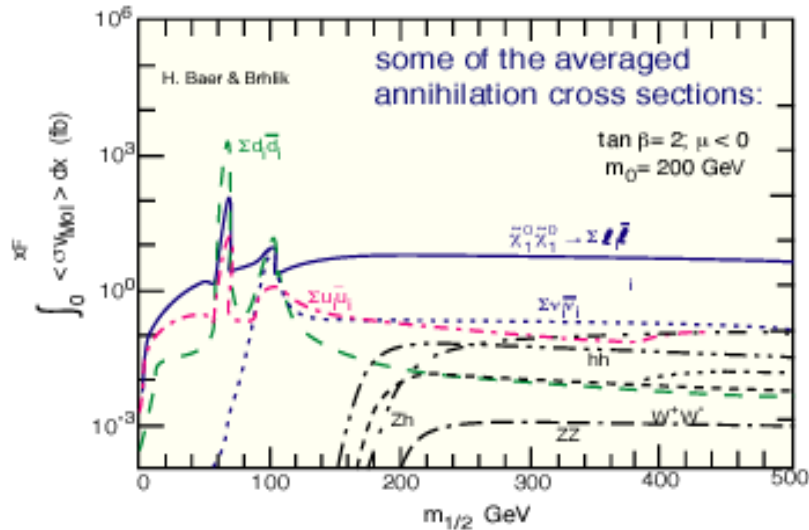


+ $\tilde{\tau}_R \tilde{\chi}_1^0$ + ...co-annihilation channels ($\tilde{\tau}_R \tilde{\chi}_1^0 \rightarrow \tau\gamma, \tau Z, \tau h$) +.....



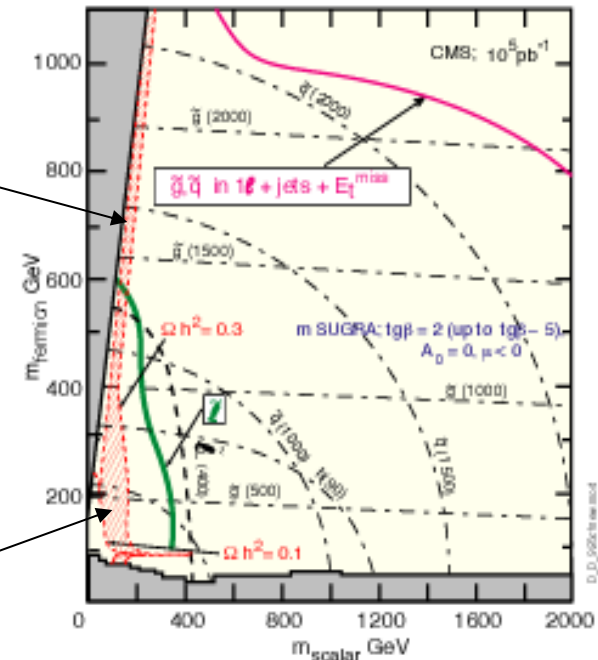
Newer Relic Densities and COBE limits in (mSUGRA) neutralino parameter space

Cross sections to calculate relic densities



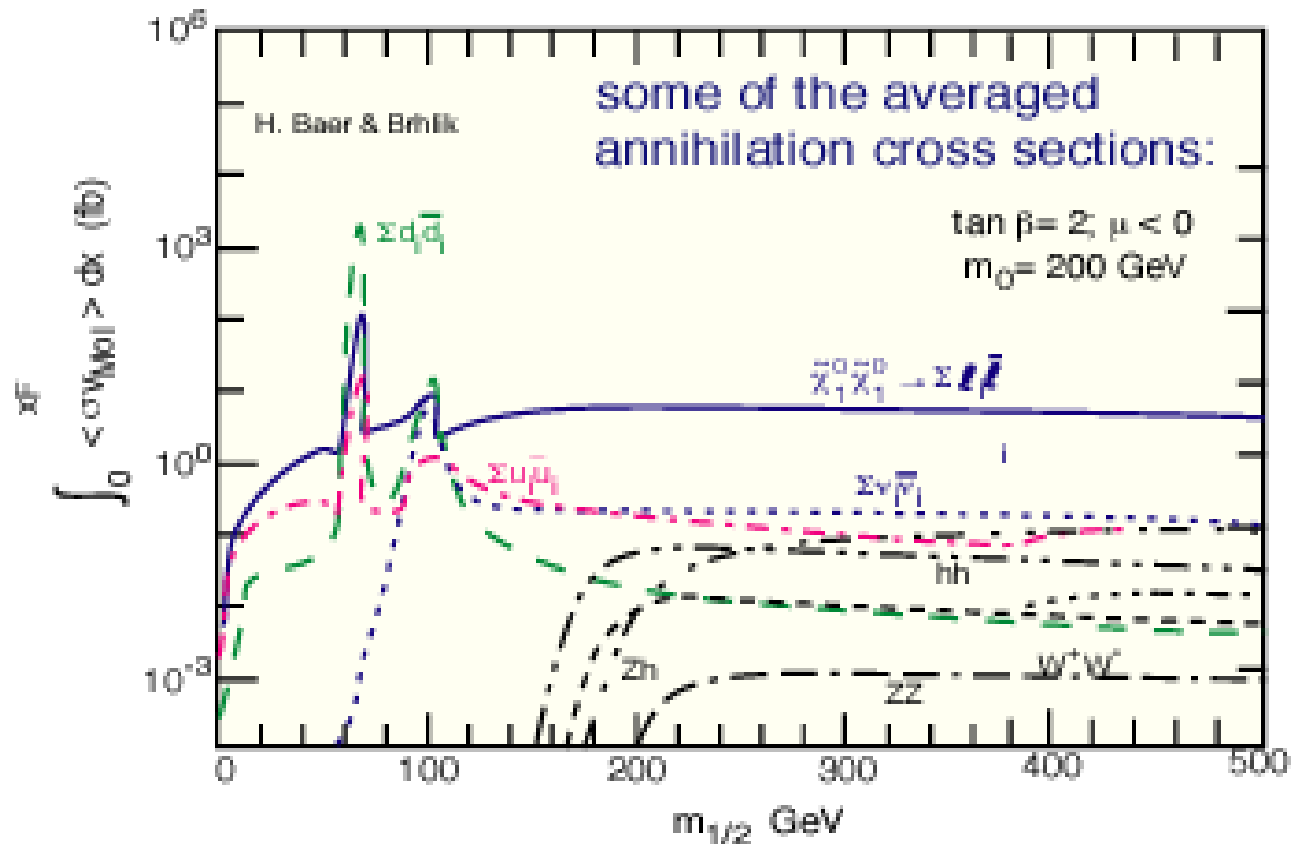
coannihilation tail

"bulk" region



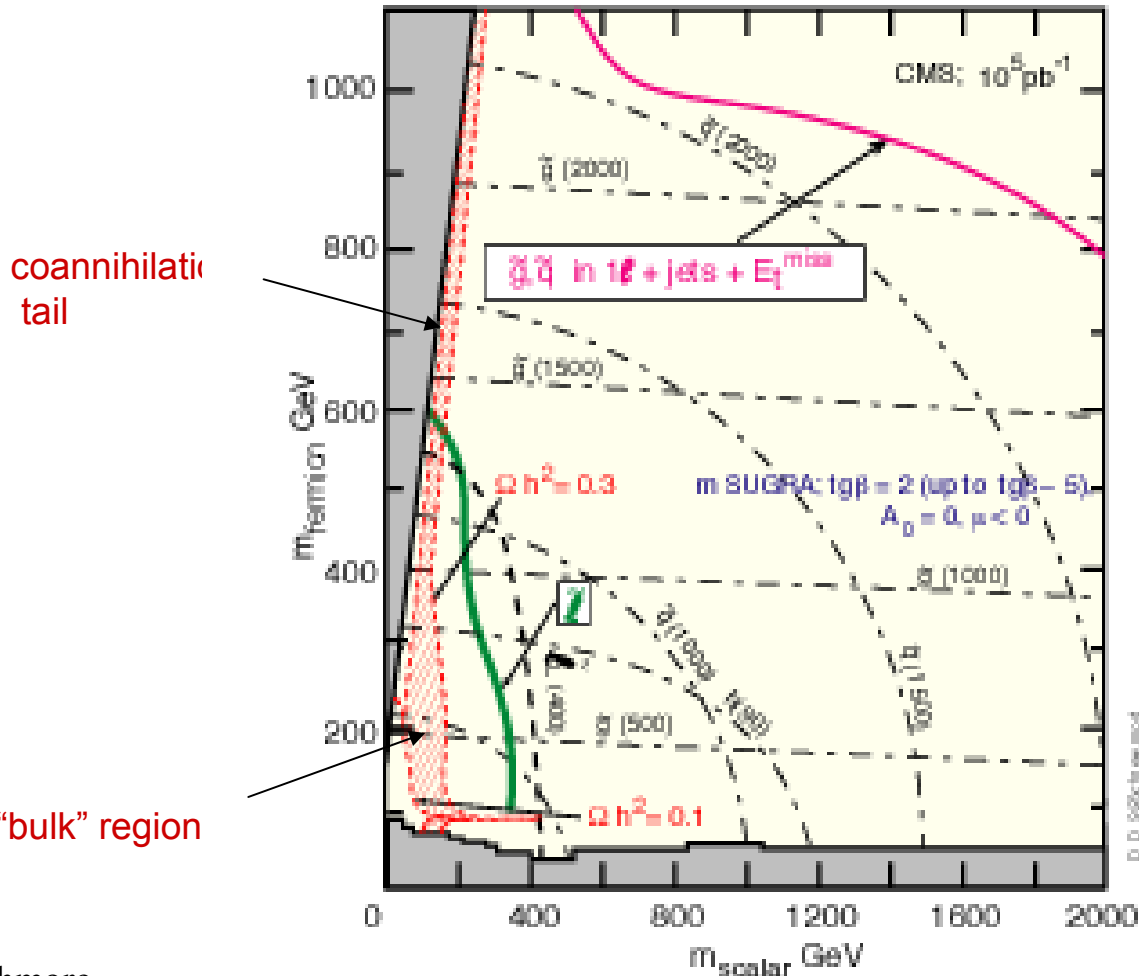
Relic densities in LHC mSUGRA parameter space (pre-WMAP)

Cross sections to calculate relic densities



Relic densities in LHC mSUGRA parameter space (pre-WMAP)

Newer Relic Densities and COBE limits in (mSUGRA) neutralino parameter space



SUSY dark matter and present constraints from laboratory experiments and WMAP

Relic neutralino DM contours, including constraints from LEP, $b \rightarrow s\gamma$, $g_{\mu-2}$ measurements and new WMAP cosmological DM constraints

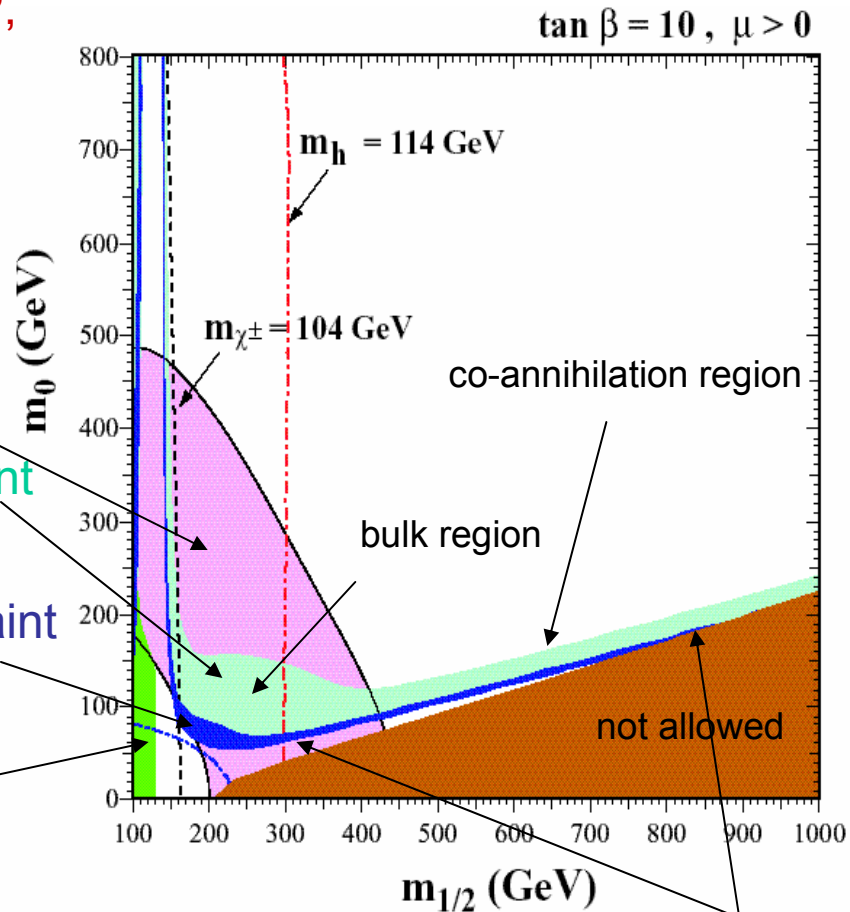
J. Ellis, K.A. Olive, Y. Santoso, V.C. Spanos, hep-ph/0303043

avored by $g_{\mu-2}$ at 2σ level

older cosmological constraint $0.1 < \Omega_{\chi} h^2 < 0.3$

newer cosmological constraint $0.094 < \Omega_{\chi} h^2 < 0.129$

excluded by $b \rightarrow s\gamma$



example at $\tan\beta = 10$

only this tail is left!

SUSY dark matter and present

J. Ellis, K.A. Olive, Y. Santoso, V.C. Spanos, hep-ph/0303043

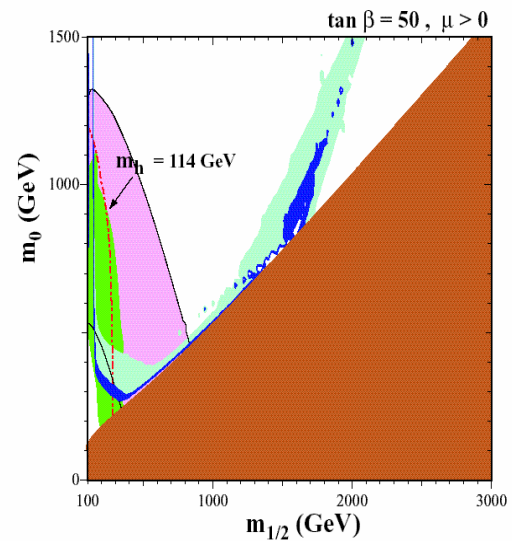
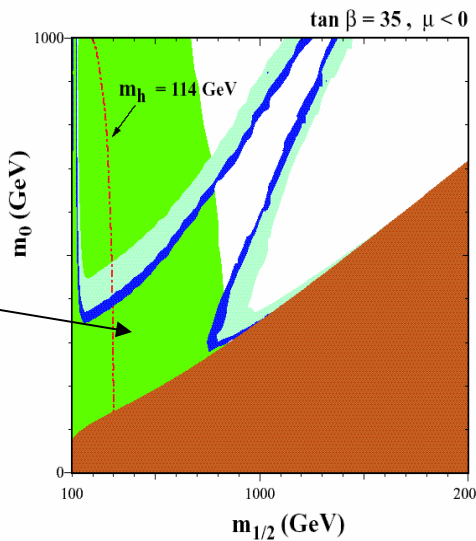
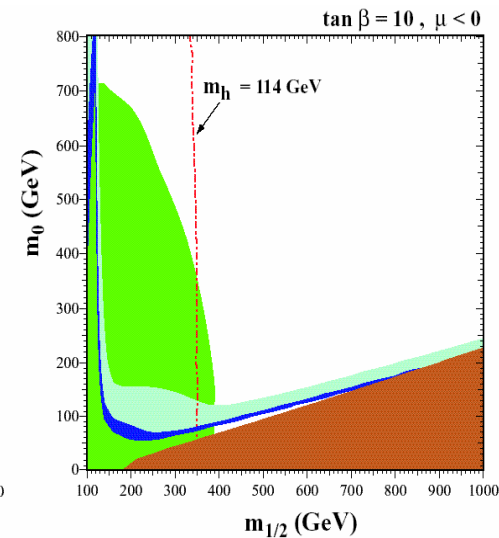
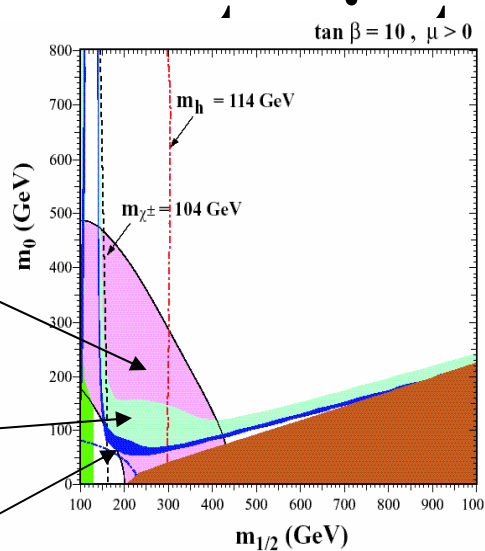
Full $\tan\beta$ range

avored by $g_{\mu-2}$
at 2σ level

older cosmological constrain
 $0.1 < \Omega_{\chi} h^2 < 0.3$

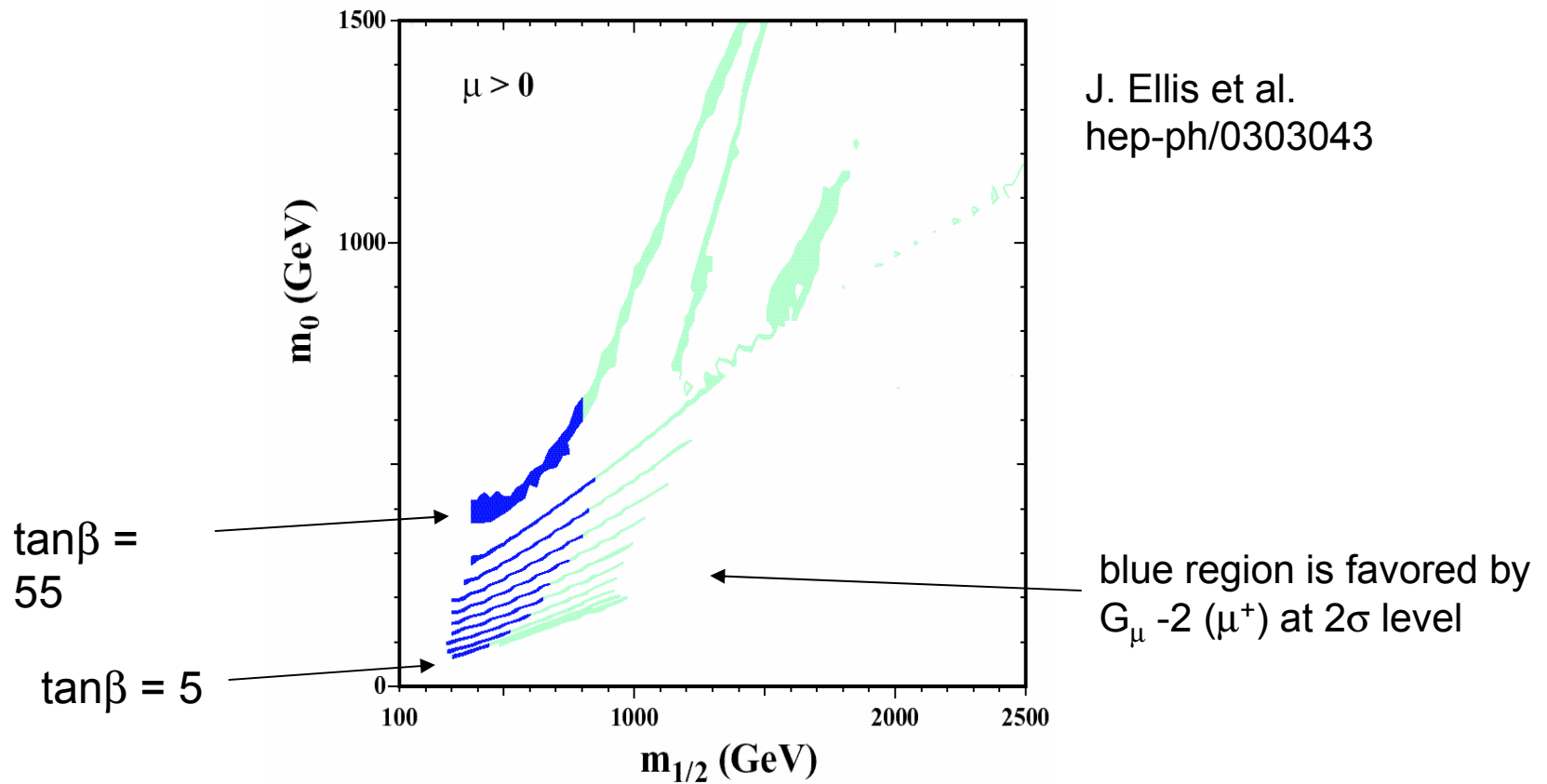
newest cosmological constra
 $0.094 < \Omega_{\chi} h^2 < 0.129$

excluded by $b \rightarrow s\gamma$

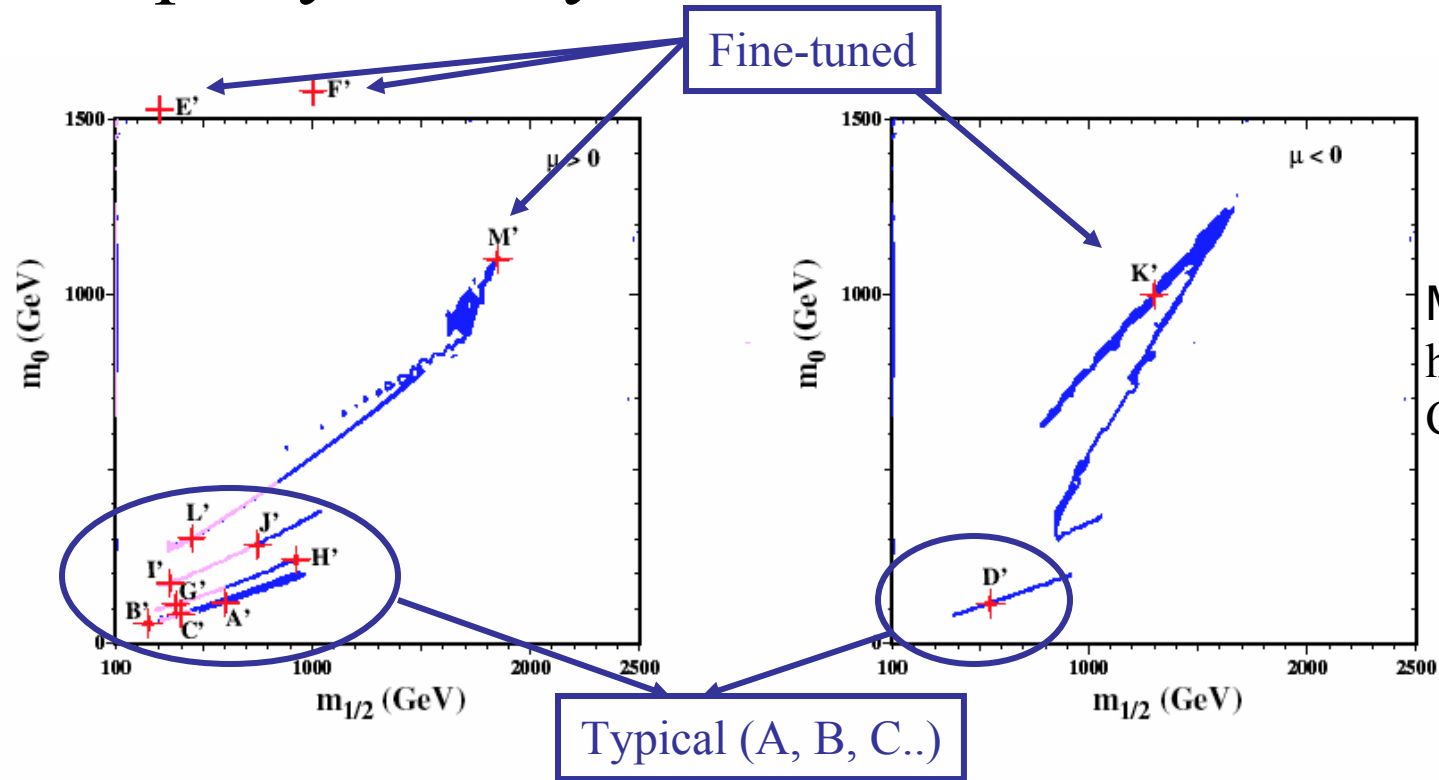


WMAP results mapped onto m_0 vs $m_{1/2}$ plane

Compatibility between WMAP results and laboratory experiments - mapping of the WMAP constraint $0.094 < \Omega_\chi h^2 < 0.129$ onto the m_0 vs $m_{1/2}$ plane for $\tan\beta$ from 5 to 55 ($\mu > 0$)



Updated Post-WMAP Benchmarks for Supersymmetry



M. Battaglia et al.
 hep-ph/0306219
 CERN-TH/2003-138

Figure 1: The shaded strips display the regions of the $(m_{1/2}, m_0)$ plane that are compatible [15] with $0.094 < \Omega_\chi h^2 < 0.129$ in the ‘bulk’, coannihilation ‘tail’, and rapid-annihilation ‘funnel’ regions, as well as the laboratory constraints, for (a) $\mu > 0$ and $\tan\beta = 5, 10, 20, 35$ and 50 , and (b) for $\mu < 0$ and $\tan\beta = 10$ and 35 . The parts of these ‘WMAP lines’ for $\mu > 0$ compatible with $g_\mu - 2$ at the $2\text{-}\sigma$ level have lighter (pink) shading [7]. The updated post-WMAP benchmark scenarios are marked in red. Points (E’,F’) in the focus-point region have larger values of m_0 .

WMAP DM results and compatibility with LHC SUSY reach

Squark, Gluino reach at the LHC
- reach in jets + E_T^{miss}

LHC/CMS reach vs SUSY DM regions after WMAP results

$$m_\chi = 0.4 m_{1/2}$$

new LEP, $b \rightarrow s\gamma$,
WMAP-compatible region
 $\Omega_{\text{CDM}} h^2 = 0.09-0.13$

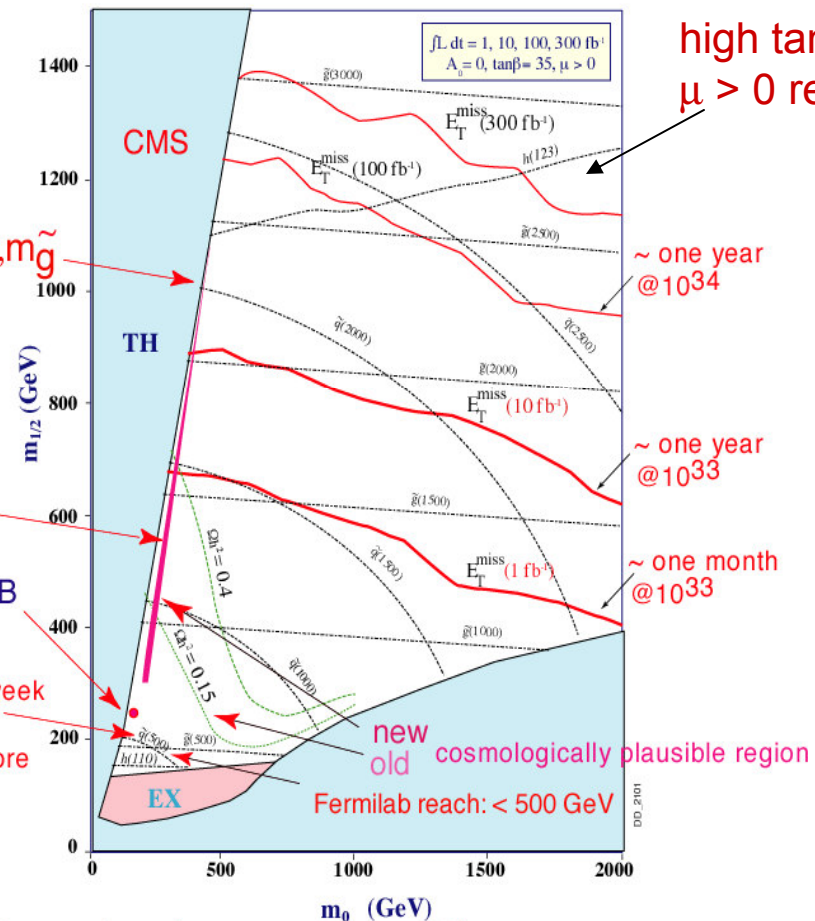
"old" Cold Dark Matter range: $\Omega_{\text{CDM}} h^2 = 0.1-0.3$

→ neutralino mass < ~ 7-800 GeV

new WMAP range: $\Omega_{\text{CDM}} h^2 = 0.09-0.13$

→ neutralino mass < ~ 4-500 GeV (for $\tan\beta < 50, \mu > 0$ or for $\tan\beta < 35, \mu < 0$)

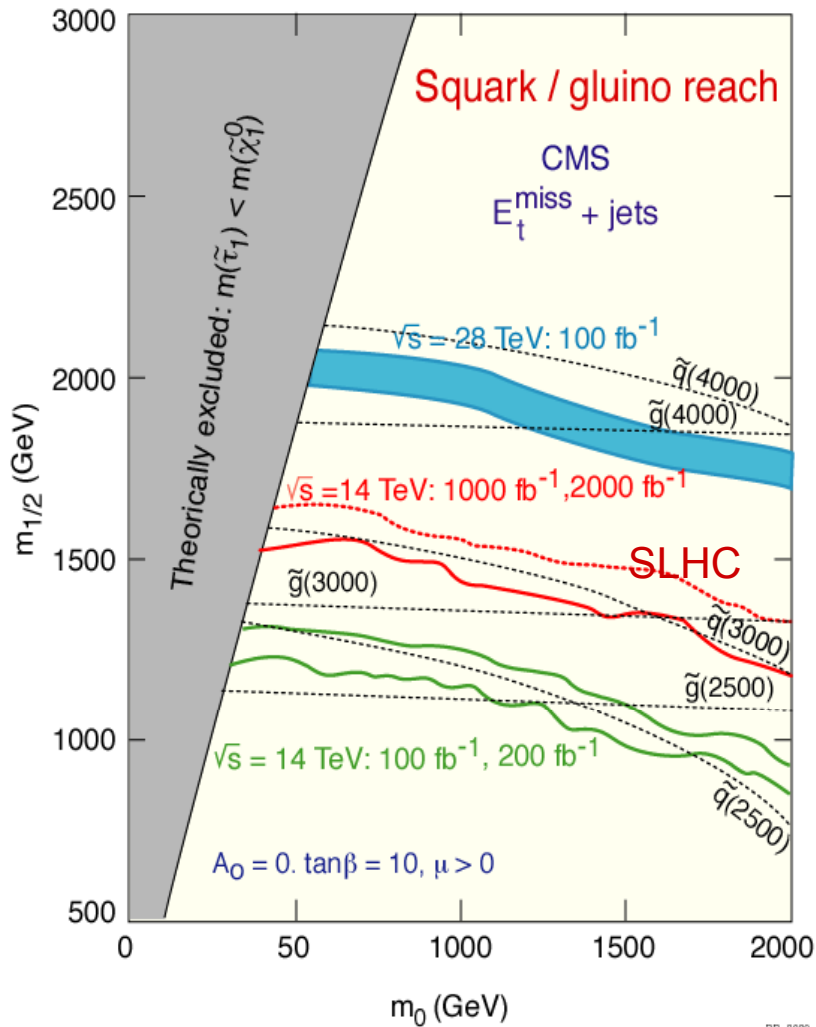
Reach vs luminosity
CMS \tilde{q}, \tilde{g} mass reach in E_T^{miss} + jets inclusive channel for various integrated luminosities



Wulcano 2004.25

D. Denegri, CMS week June 2003.04

SUSY at SLHC (I)



Higher integrated luminosity brings an obvious **increase in mass reach** in squark, gluino searches, i.e. **in SUSY discovery potential**; this is not too demanding on detectors as very high E_t jets, E_t^{miss} are involved, large pile-up not so detrimental

⇒ with SLHC the **SUSY reach** is increased by ~ 500 GeV, **up to ~ 3 TeV** in squark and gluino masses

but this is just “the reach”, **the main advantage of increased statistics should be in the sparticle spectrum reconstruction possibilities**, larger fraction of spectrum, more precision, but this would require detectors of comparable performance to “present ones”

Notice advantage of a 28 TeV machine.

Model dependent tests, sparticle masses, neutralino mass

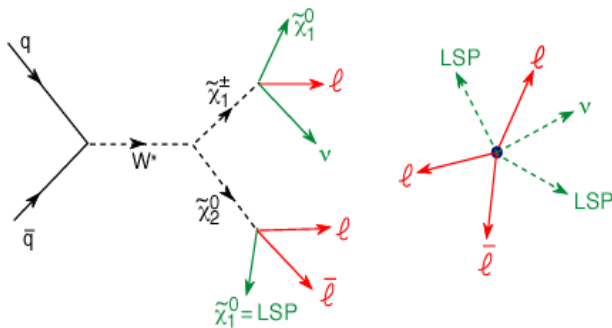
- If a viable DM candidate is found initially (events with large missing E_t , stable) assume specific consistent model e.g. CMSSM / mSUGRA.
- Measure model parameters (m_0 , $m_{1/2}$, $\tan(\beta)$, $\text{sign}(\mu)$, A_0 for mSUGRA).
- Check consistency with accelerator constraints (m_h , $g_{\mu-2}$, $b \rightarrow s\gamma$ etc.)
- Estimate $\Omega_\chi h^2 \rightarrow$ consistency check with astrophysics (WMAP etc.)
- Ultimate test of DM only possible in conjunction with astroparticle experiments

\rightarrow measure m_χ , $\sigma_{\chi p}$, ϕ_{sun} etc.

neutralino mass in exclusive chargino-neutralino final states

How to look for charginos and neutralinos at LHC

- Production, for ex by Drell Yan:



➔ Search in: $3\ell^\pm$ and no jets + (E_T^{miss}) events

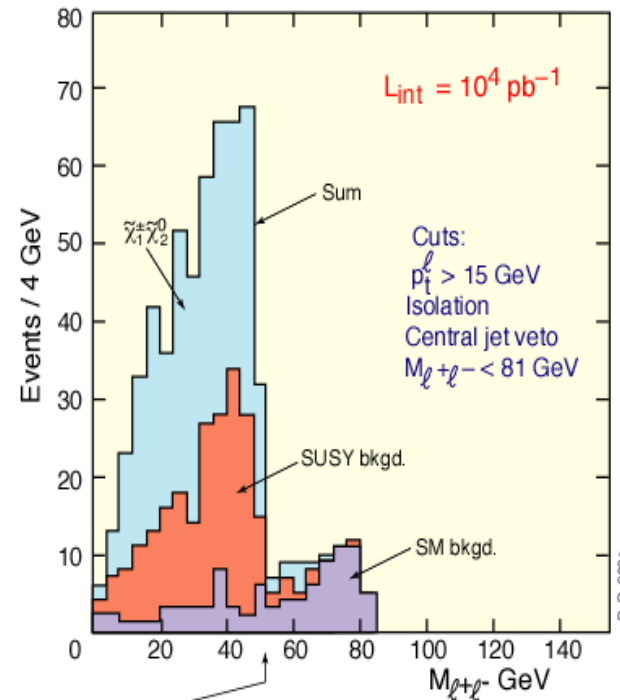
Require:

- Three isolated leptons: $p_T^\ell > 15$ GeV
- Veto central jets with $E_T > 25$ GeV in $|\eta^j| < 3.5$
- $m_{\ell\bar{\ell}} < 81$ GeV or $m_{\ell\bar{\ell}} \neq m_z \pm \delta m_z$
- $E_T^{\text{miss}} < \dots$ should be significant

Backgrounds: $t\bar{t}$, WZ, ZZ, Zbb, $b\bar{b}$ other SUSY channels ($\tilde{g}, \tilde{q}, \tilde{\ell}, \chi^0, \chi^\pm$)

$\tilde{\chi}_1^0$ mass determination in $3\ell^\pm$ + no jets + E_T^{miss} final state from $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ production

$m_0 = 200$ GeV, $m_{1/2} = 100$ GeV, $\tan\beta = 2$, $A_0 = 0$, $\mu < 0$
 $M(\tilde{\chi}_2^0) - M(\tilde{\chi}_1^0) \approx M(\tilde{\chi}_1^0) \approx 52$ GeV



$$M_{\ell\ell}^{\text{max}} = M(\tilde{\chi}_2^0) - M(\tilde{\chi}_1^0) \approx M(\tilde{\chi}_1^0) \rightarrow m_{1/2} \approx 2.5 M(\tilde{\chi}_1^0)$$

in m SUGRA

R.Cashmore the "edge" can be measured with better than ~ 1 GeV precision and thus neutralino masses within this model - and where the edge is visible! Dark Matter 3 39

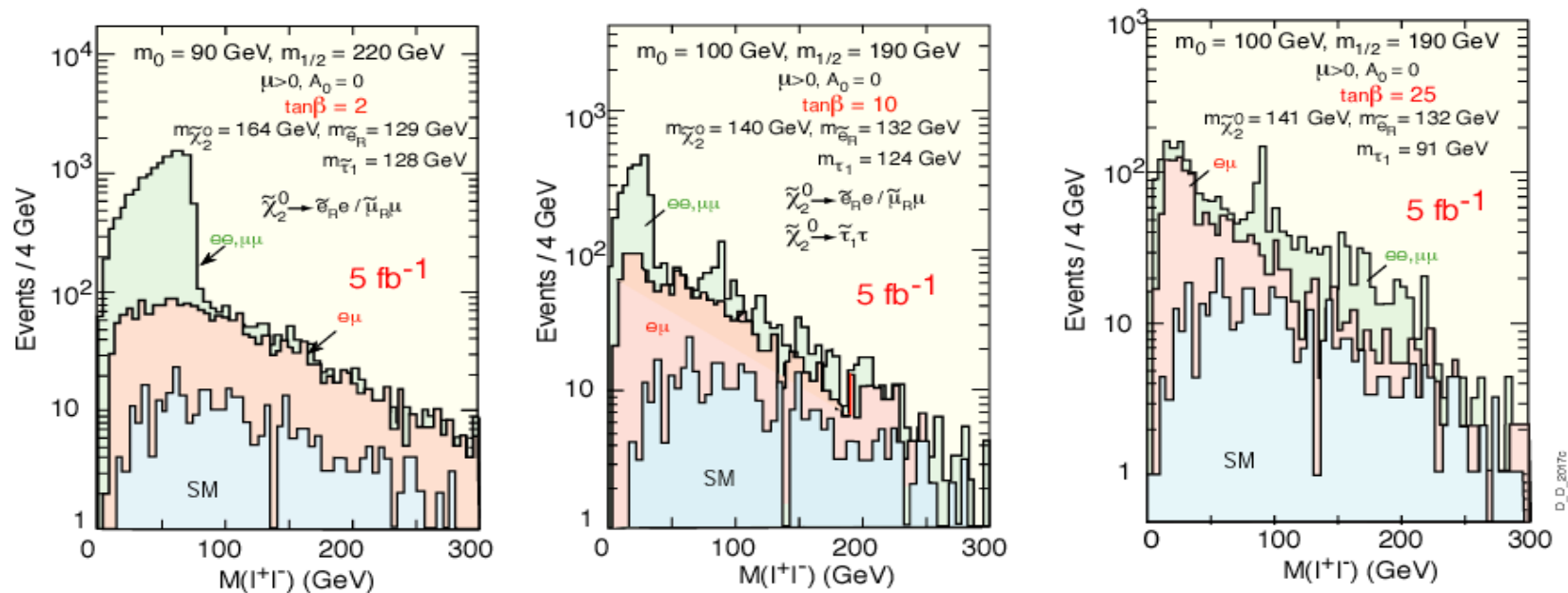
Evolution of inclusive dilepton spectra with increasing $\tan\beta$

Inclusive dilepton edges

($m_0 \sim 100$ GeV, $m_{1/2} \sim 200$ GeV)

5 fb⁻¹

inclusive isolated opposite-sign dilepton spectra (ee, $\mu\mu$, e μ)



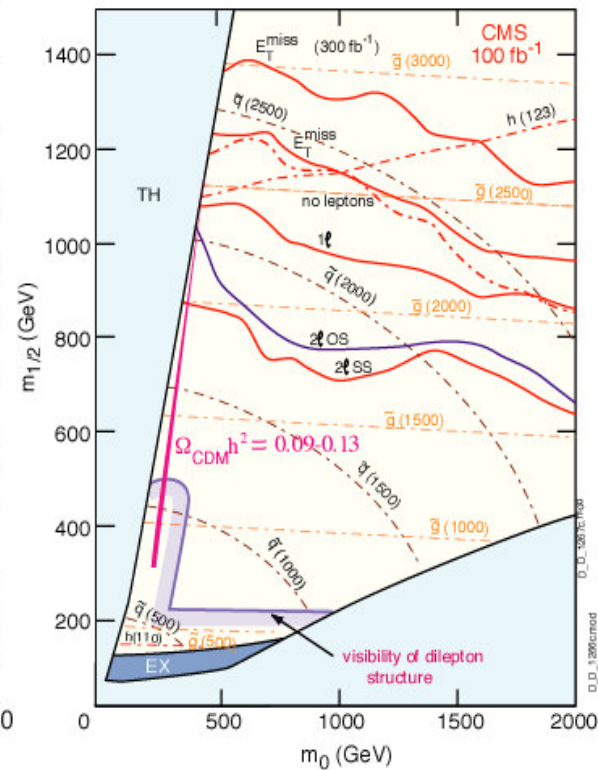
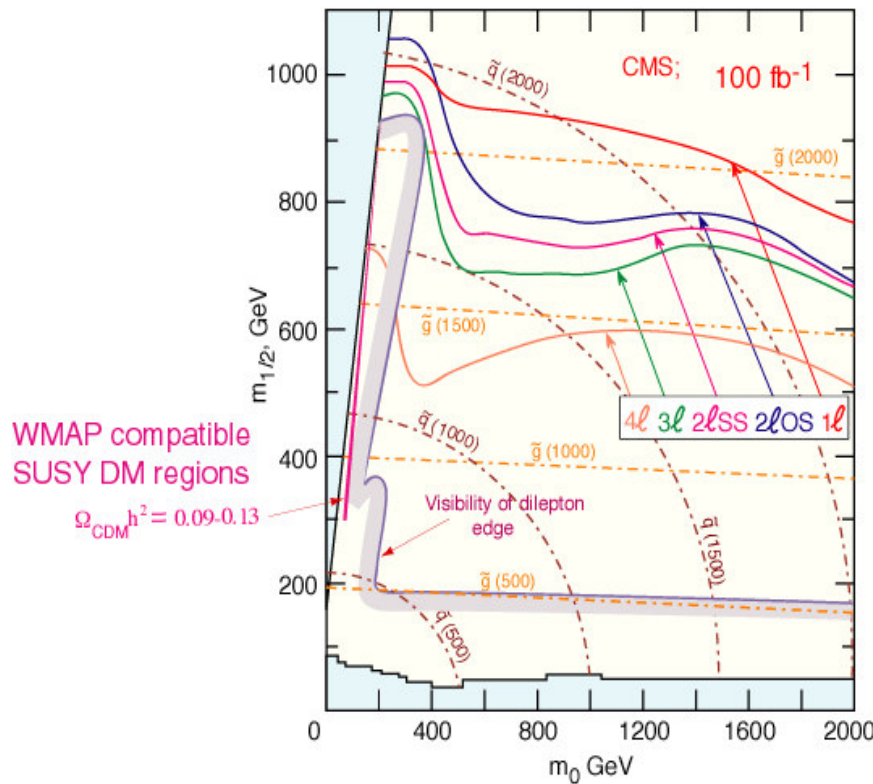
→ SUSY may reveal itself early through peculiarities in dilepton spectra, for example such "edges" or "collapsed edges"

WMAP- DM contours and dilepton structures

Observability of " $l+l^-$ edge" in inclusive isolated two - leptons + E_t^{miss} final states

m SUGRA; $\tan \beta = 2$, $A_0 = 0$, $\mu < 0$
all leptons isolated

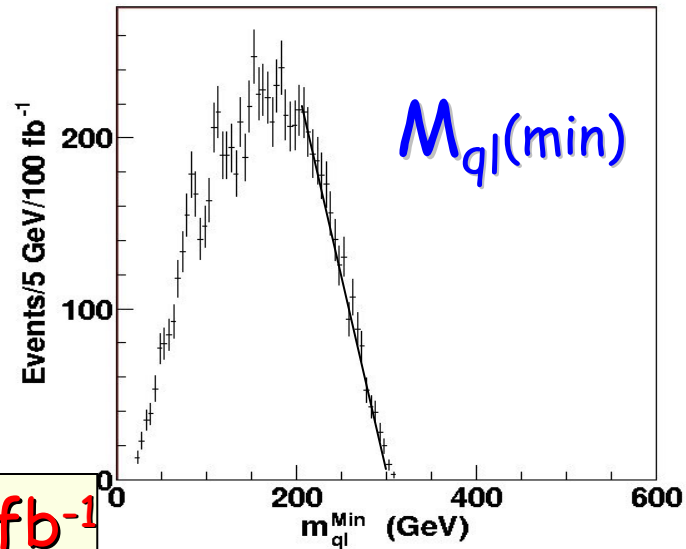
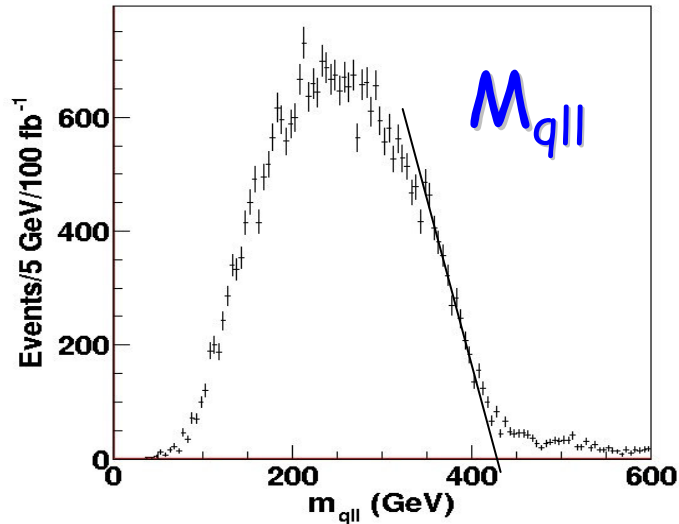
m SUGRA, $A_0 = 0$, $\tan \beta = 35$, $\mu > 0$
5 σ contours ; non - isolated muons



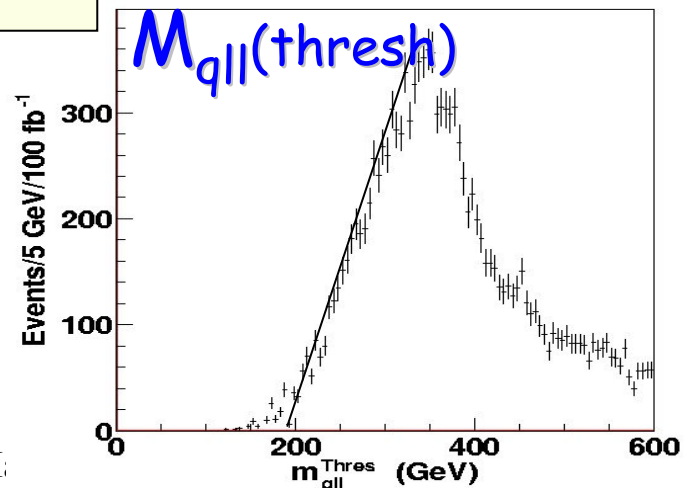
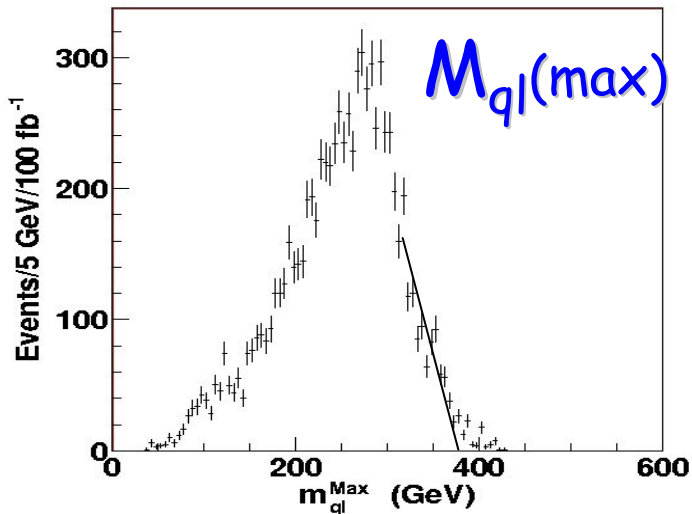
high $\tan \beta$ regime

Sparticle reconstruction, use several edges

Post-LEP SUSY benchmark point B, ATLAS/2004-007



100 fb⁻¹



Dark M.

SUMMARY

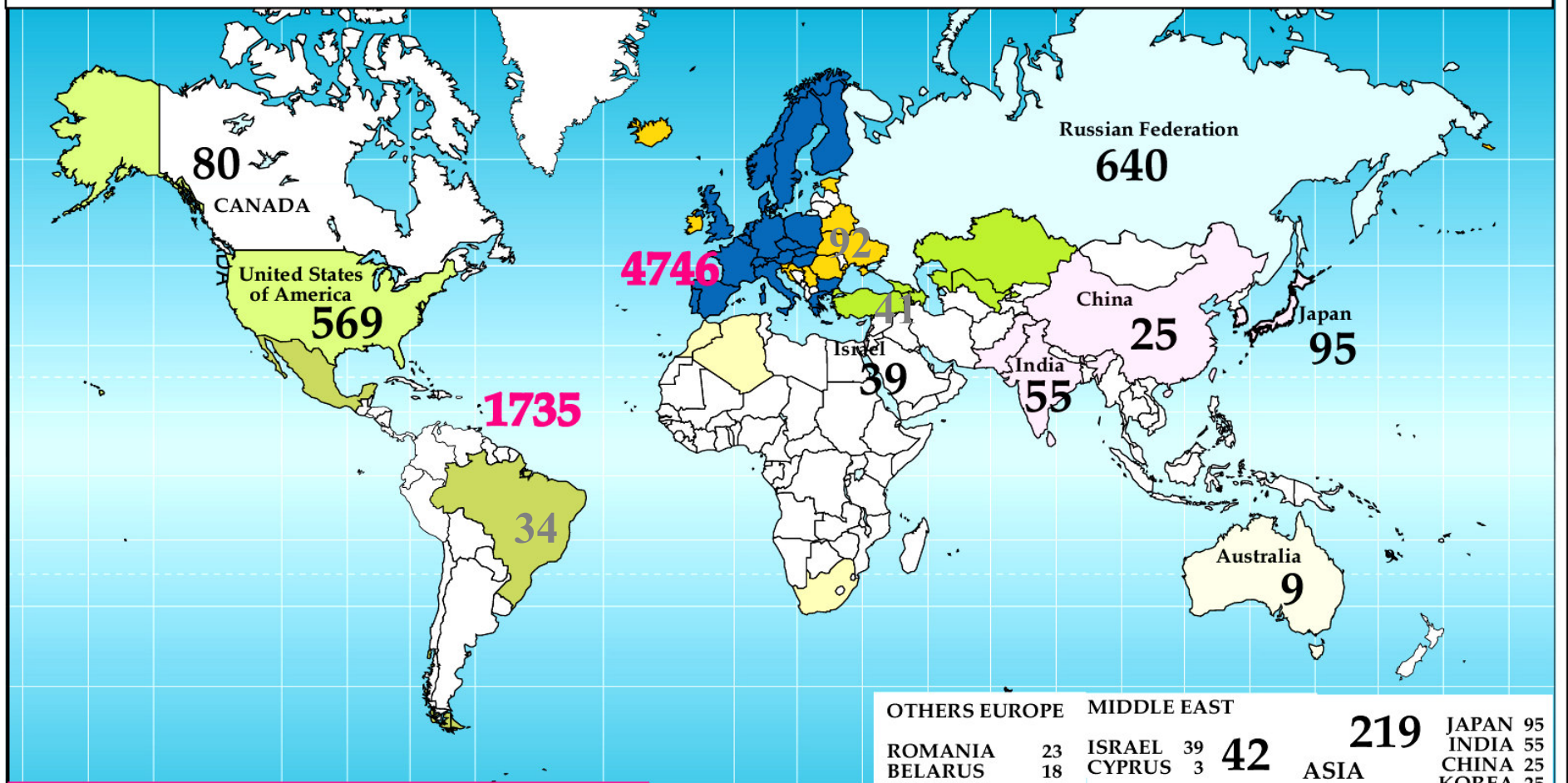
- LHC and Direct Searches will allow an understanding of DM

The next 4 years should be an exciting time

Particular Thanks

- Alex Murphy and Tim Sumner
- Colleagues at the Gran Sasso
- Daniel Denegri
and
- Numerous theorists

Worldwide Scientific Collaboration



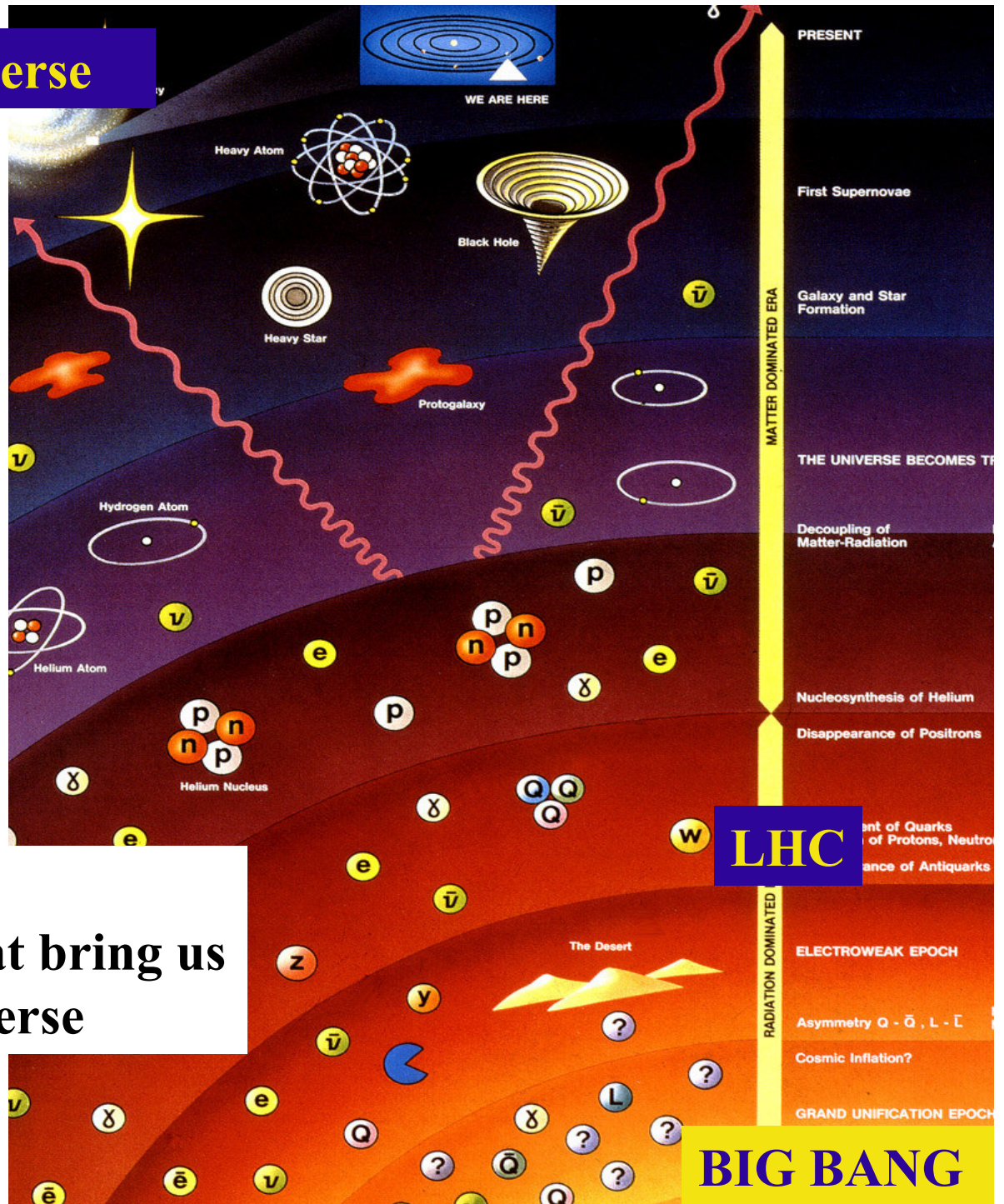
MEMBER STATES		
AUSTRIA	GERMANY	4746
BELGIUM	GREECE	
BULGARIA	HUNGARY	
CZECH REPUBLIC	ITALY	
DENMARK	NETHERLANDS	
FINLAND	NORWAY	
FRANCE	POLAND	
	PORTUGAL	
	SLOVAKIA	
	SPAIN	
	SWEDEN	
	SWITZERLAND	
	UNITED KINGDOM	

LATIN AMERICA	
BRAZIL	26
MEXICO	8
TOTAL	34

OTHERS EUROPE	MIDDLE EAST	ASIA
ROMANIA 23	ISRAEL 39	219
BELARUS 18	CYPRUS 3	
UKRAINE 17		
SLOVENIA 15	TOTAL 42	TOTAL 219
CROATIA 12		
SERBIA 2		
IRELAND 2		
ESTONIA 2		
MALTA 1		
	TOTAL 9	
	MOROCCO 4	41
	SOUTH AFRICA 4	
	ALGERIA 1	
		TURKEY 17
		ARMENIA 15
		GEORGIA 7
		AZERBAIJAN 1
		UZBEKISTAN 1

The History of the Universe

Time ↑



Particle Accelerators
are Time-Machines that bring us
back to the Early Universe

R.Cashmore

Perspectives for Dark Matter Searches in ATLAS

Dan Tovey

University of Sheffield

SUSY Dark Matter studies at ATLAS will proceed in four stages:



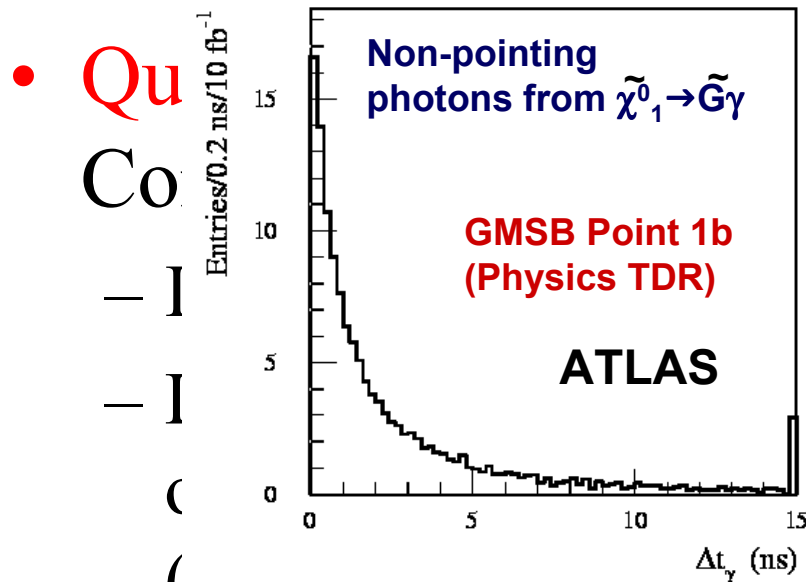
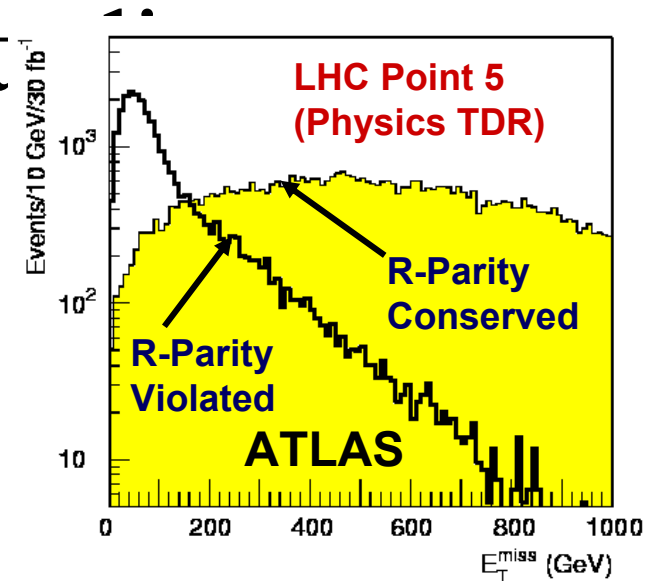
1) SUSY Discovery phase (discussed by Marco for CMS)

→ success assumed!

2) Inclusive Studies (measurement of SUSY Mass Scale, comparison of significance in inclusive channels).
In this talk focus on Stages 2, 3 and 4

3) Exclusive studies and interpretation within specific model framework (e.g. Constrained

- Following any discovery of **Inclusive St** SUSY next task will be to test broad features of potential Dark Matter candidate.



- Question 2: Is the LSP the lightest neutralino?**

 - Natural in many MSSM models
 - If YES then test for consistency with astrophysics
 - If NO then what is it?
 - e.g. Light Gravitino DM from GMSB models (not considered here)

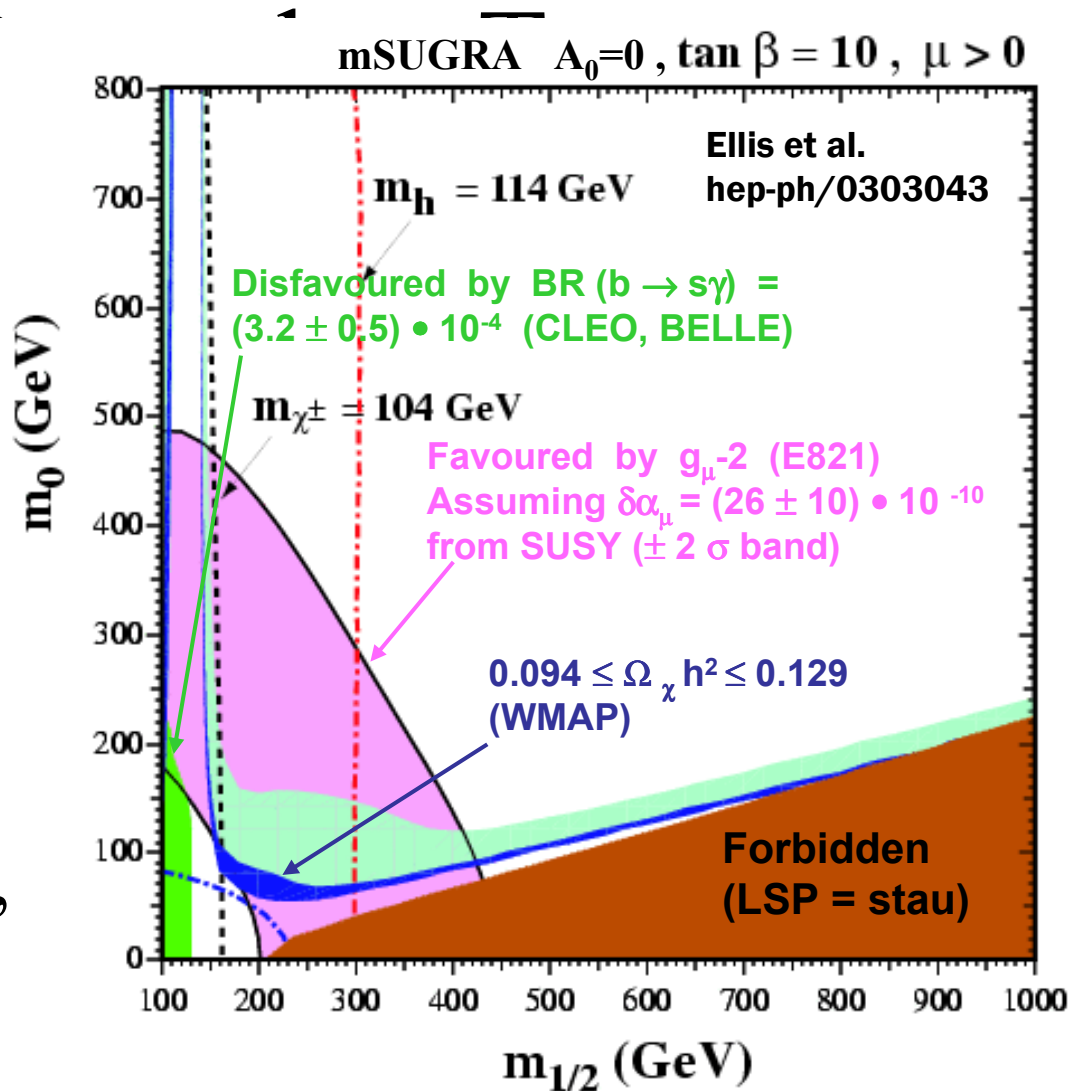
- If a viable **Model-C** candidate is found initially assume specific consistent model

– e.g. CMSSM / mSUGRA.

- Measure model parameters (m_0 , $m_{1/2}$, $\tan(\beta)$, $\text{sign}(\mu)$, A_0 for CMSSM).

R.Cashmore

- Check consistency



Dark Matter 3

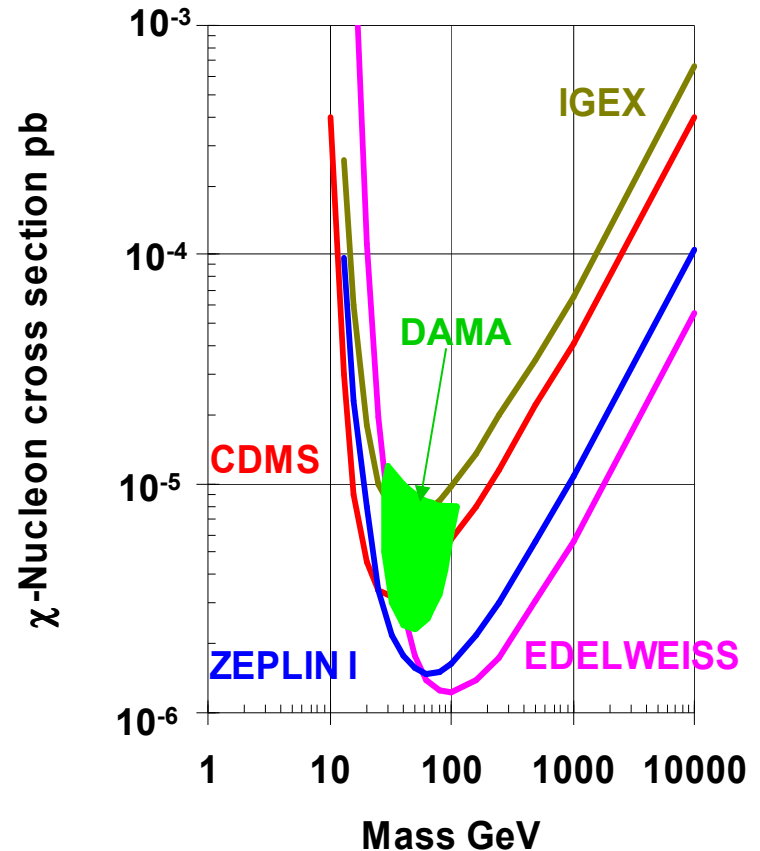
50

Terrestrial Dark Matter

- Terrestrial Dark Matter Experiments

- Direct searches for elastic scattering of Dark Matter particles from atomic nuclei ($\sigma_{\chi p}$)
- Indirect searches via self-annihilation products from e.g. centre of sun (e.g. high energy neutrino flux ϕ_{sun})

- Direct search experiments can measure neutralino



Current world-best limits (ZEPLIN-I, EDELWEISS):
 $\sigma_{\chi p} \lesssim 10^{-6} \text{ pb}$ ($m_{\chi} \sim 100 \text{ GeV}$)

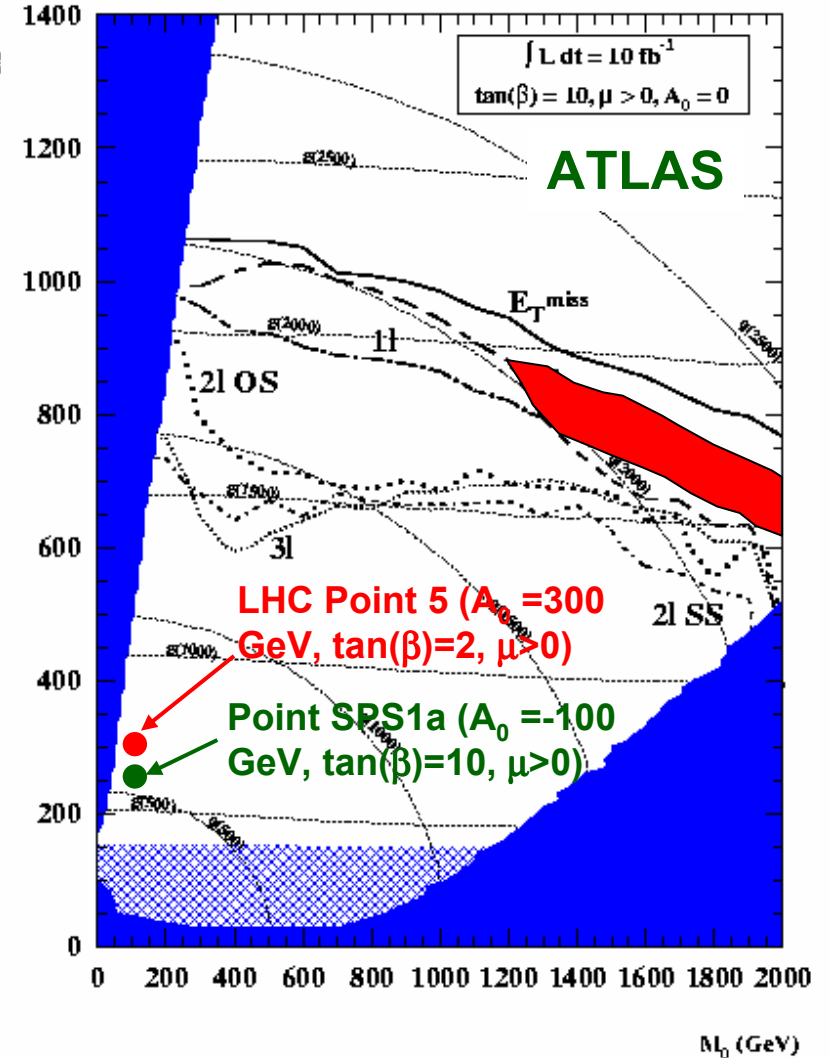
Measuring CMSSM Parameters

- First indication of CMSSM parameters from inclusive channels

– Compare significance in jets + E_T^{miss} + n leptons channels

Point	m_0	$m_{1/2}$	A_0	$\tan(\beta)$	$\text{sign}(\mu)$
LHC Point 5	100	300	300	2	+1
SPS1a	100	250	-100	10	+1

Sparticle	Mass (LHC Point 5)	Mass (SPS1a)
\tilde{q}_L	~690 GeV	~530 GeV
$\tilde{\chi}_2^0$	233 GeV	177 GeV
\tilde{l}_R	157 GeV	143 GeV
$\tilde{\chi}_1^0$	122 GeV	96 GeV



previously:

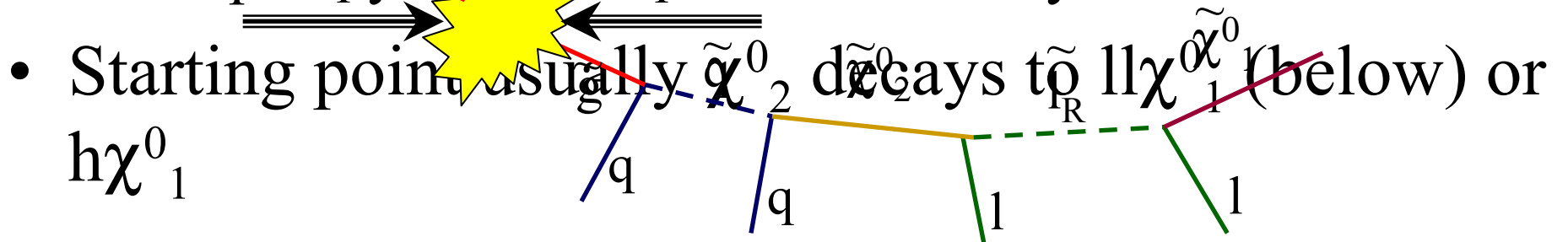
K. Cashmore

Dark Matter 3

52

- Fit reconstructed mass combinations obtain parameters

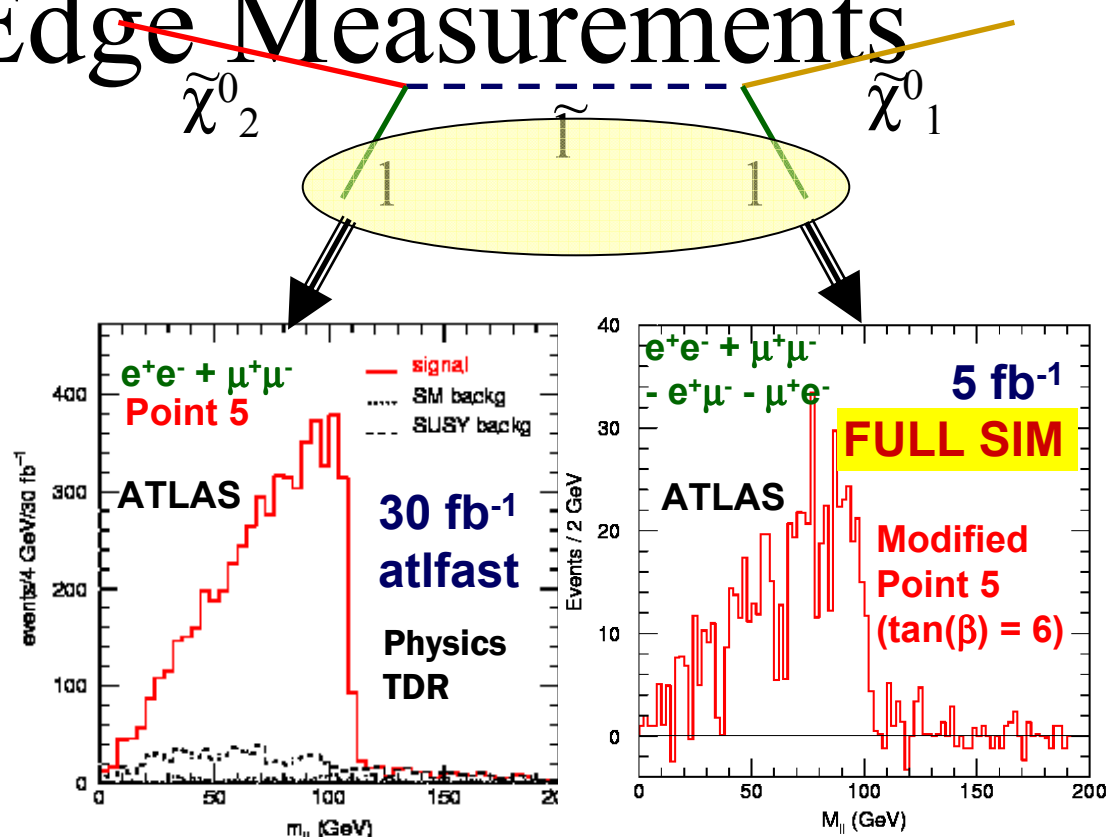
– Helped by low SUSY cascade decay chains



- Two neutral LSPs escape from each event

Dilepton Edge Measurements

- When kinematically accessible $\tilde{\chi}_2^0$ can undergo sequential two-body decay to $\tilde{\chi}_1^0$ via a right-slepton (e.g. LHC Point 5).

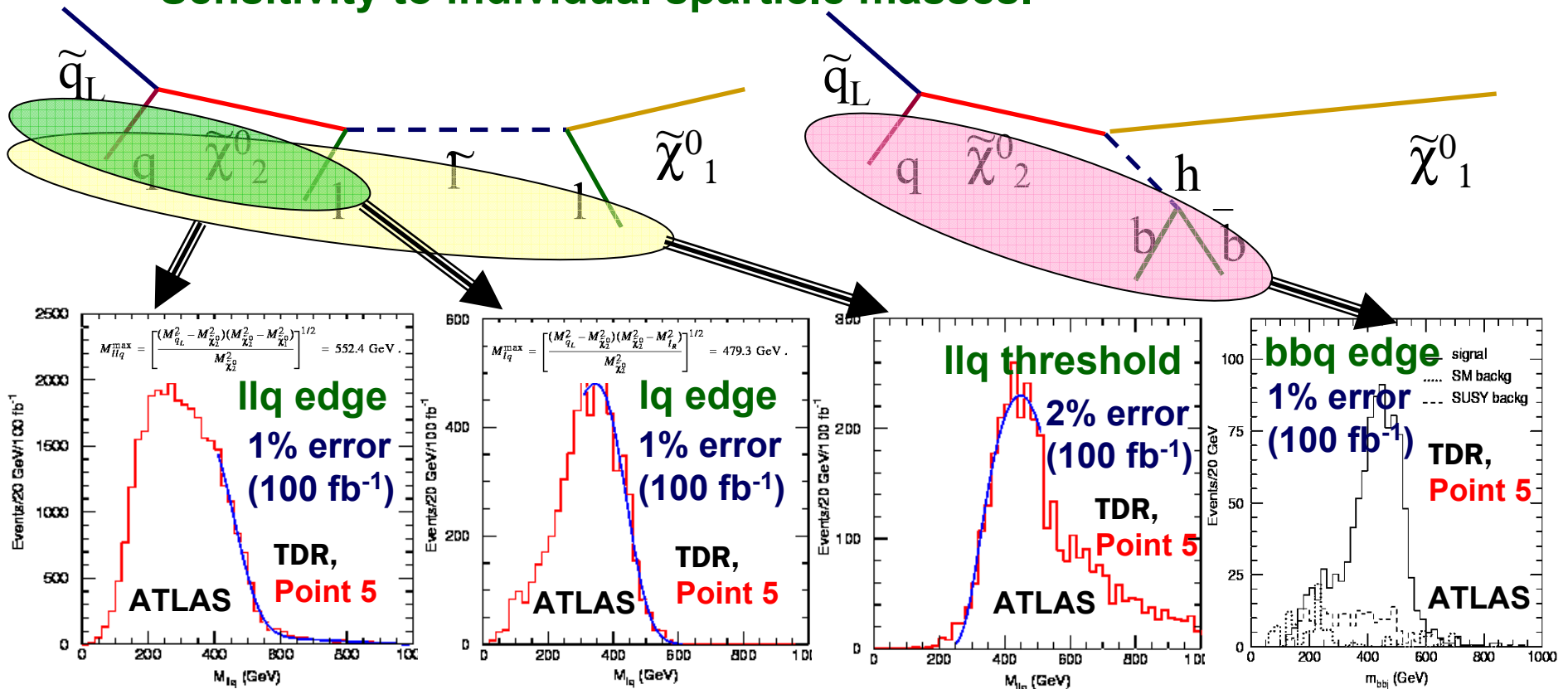


- Results in sharp OS SF dilepton invariant mass edge sensitive to

$$M_{ll}^{\max} = M(\tilde{\chi}_2^0) \sqrt{1 - \frac{M^2(\tilde{l}_R)}{M^2(\tilde{\chi}_2^0)}} \sqrt{1 - \frac{M^2(\tilde{\chi}_1^0)}{M^2(\tilde{l}_R)}} = 108.93 \text{ GeV}$$

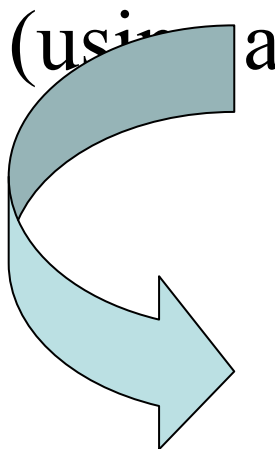
- Dilepton edge starting point for reconstruction of decay chain.
- Make invariant mass combinations of leptons and jets.
- Gives multiple constraints on combinations of four masses.
- Sensitivity to individual particle masses.

Measurements Involving Squarks

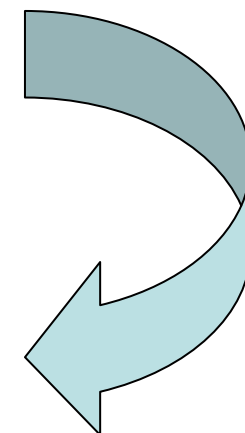


- # Measuring CMSSM Parameters

Within CMSSM a most direct approach is to calculate edge positions using sparticle masses and formulae

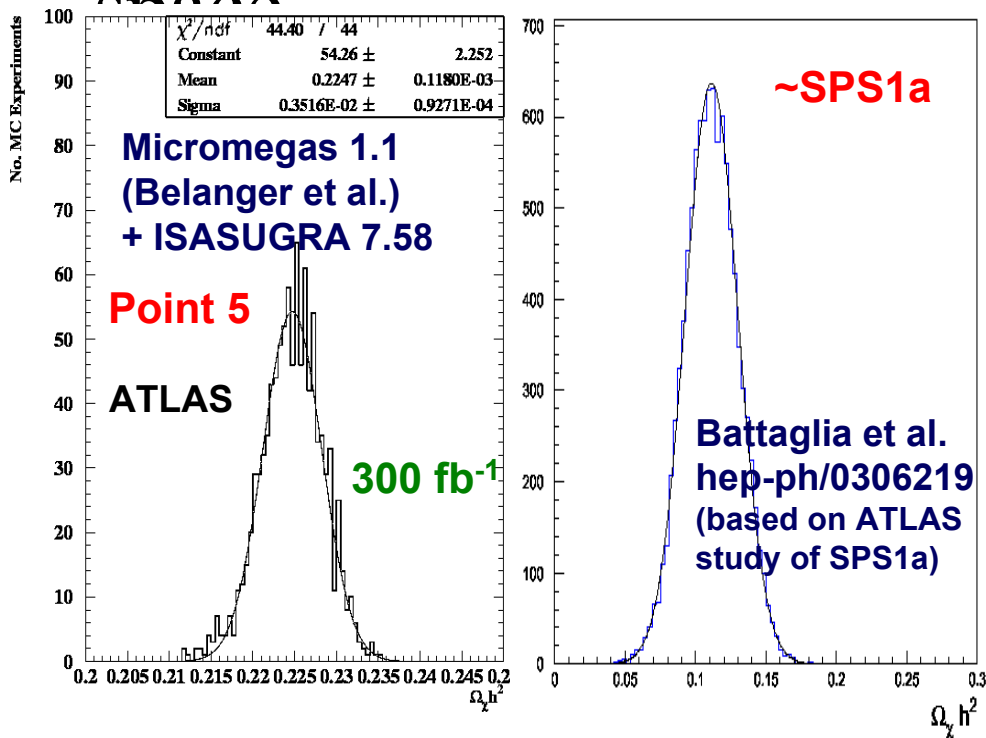


Point	m_0	$m_{1/2}$	A_0	$\tan(\beta)$	$\text{sign}(\mu)$
LHC Point 5	100	300	300	2	+1
SPS1a	100	250	-100	10	+1

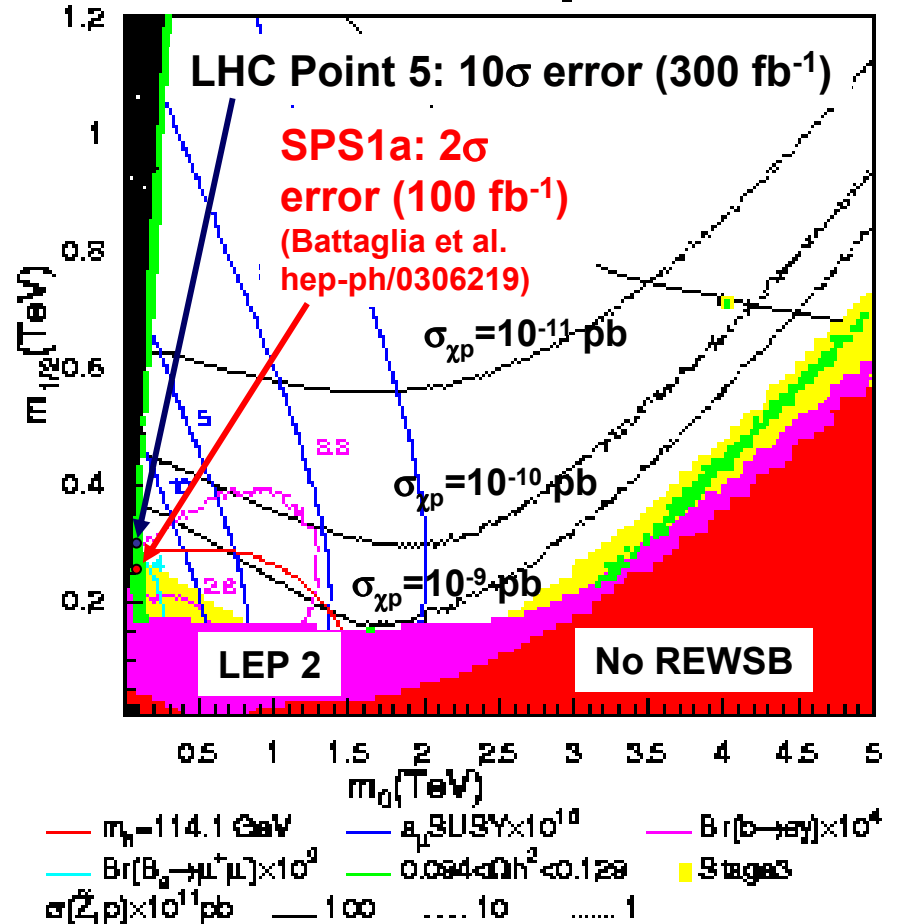


Parameter	Expected precision	
	30 fb^{-1}	300 fb^{-1}
m_0	$\pm 3.2\%$	$\pm 1.4\%$
$m_{1/2}$	$\pm 0.9\%$	$\pm 0.6\%$
$\tan(\beta)$	$\pm 0.5\%$	$\pm 0.5\%$

- Measurement precision **Relic Density** Baer et al. hep-ph/0305191
 $mSUGRA, \tan \beta=10, A_0=0, \mu>0$
 for $\Omega_\chi h^2$ depends strongly on region of parameter



R.Cashmore

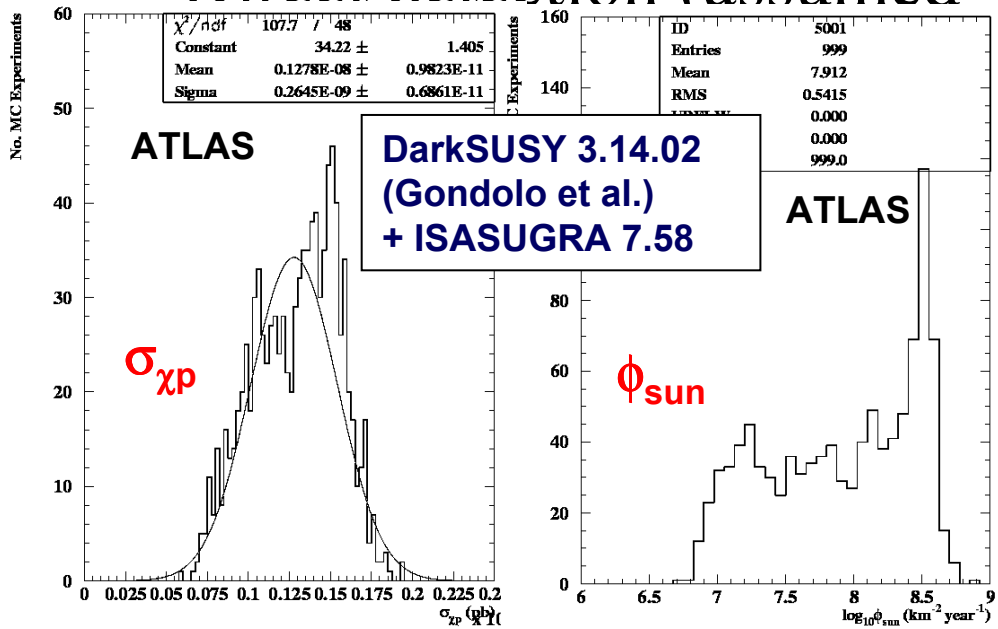


Dark Matter 3

- Also use **Dark Matter Sensitivity** parameters to predict signals observed in terrestrial dark matter searches (Point 5)

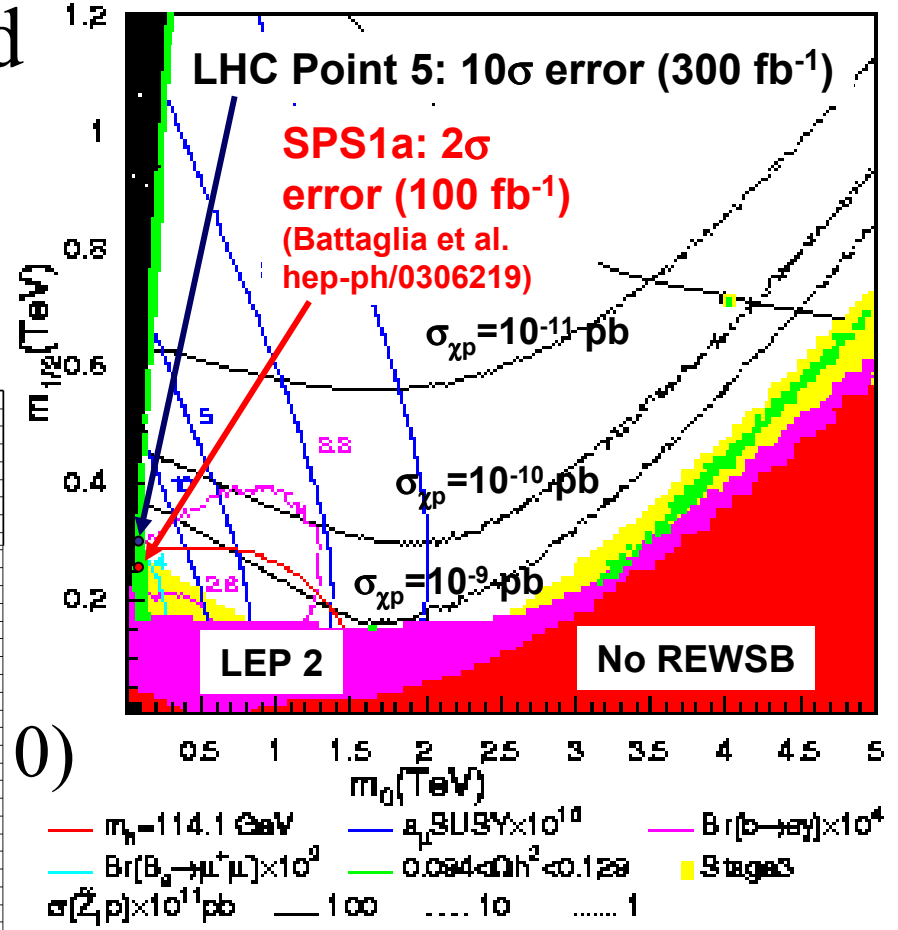
Baer et al. hep-ph/0305191
 mSUGRA, $\tan\beta=10, A_0=0, \mu>0$

– Direct detection (assumed)

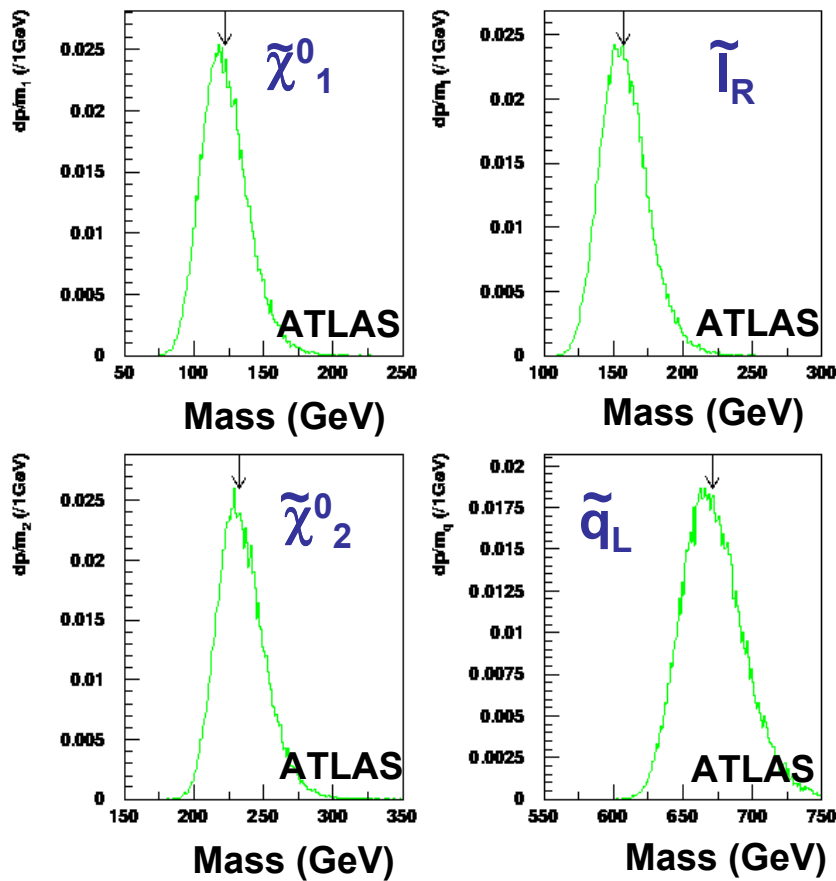


R.Cashmore

Dark Matter 3



- Alternative approach to CMS SM fit to edge positions.



f simultaneous

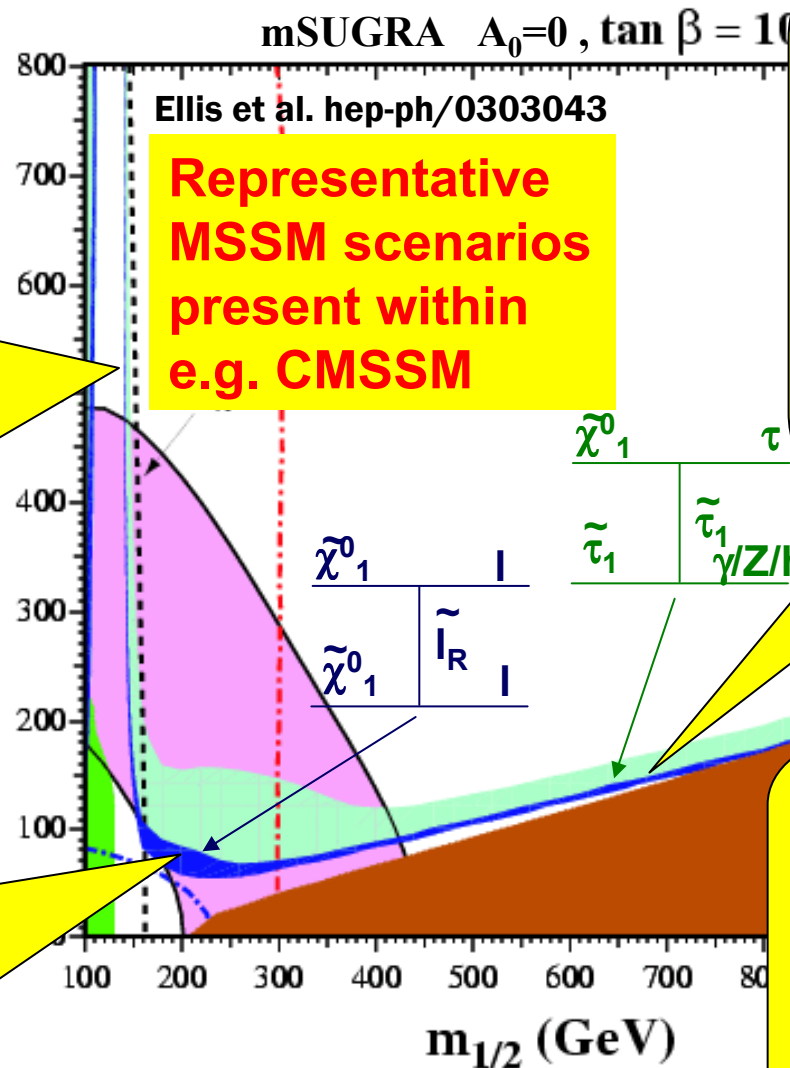
Similar process for $\tilde{\tau}_1$ mass at high $\tan(\beta)$

Sparticle	Expected precision (100 fb ⁻¹)
\tilde{q}_L	$\pm 3\%$
$\tilde{\chi}^0_2$	$\pm 6\%$
\tilde{l}_R	$\pm 9\%$
$\tilde{\chi}^0_1$	$\pm 12\%$

- Use approximations together with other measurements to obtain 'model-independent' estimates of $\Omega_\chi h^2$, $\sigma_{\chi p}$, ϕ_{sun} etc.

'Focus point' region (significant \tilde{h} component to LSP): v. difficult, need $m(\tilde{\chi}^0_1)$, μ , m_A , $\tan(\beta)$ etc. + $m(t)$ to high precision. More study needed

'Bulk' region (t-channel slepton exchange - LSP mostly Bino): need $m(\tilde{\chi}^0_1)$, $m(\tilde{l}_R)$, $m(\tilde{\tau}_1)$. 'Bread and Butter' region for LHC Expts.



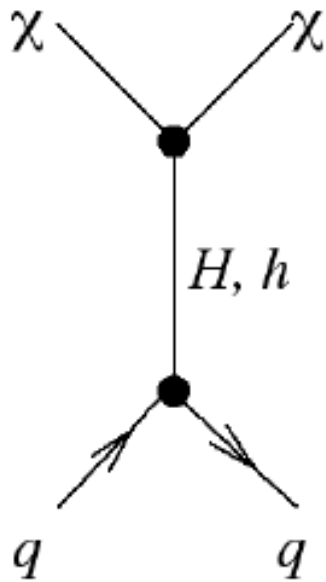
Slepton Co-annihilation region (LSP ~ pure Bino): need $m(\tilde{\chi}^0_1)$, $m(\tilde{\tau}_1)$. Small mass difference makes measurement difficult however.

Also 'rapid annihilation funnel' at Higgs pole at high $\tan(\beta)$: $m(\tilde{\chi}^0_1)$, m_A , μ , $\tan(\beta)$, $m(t)$ etc. needed.

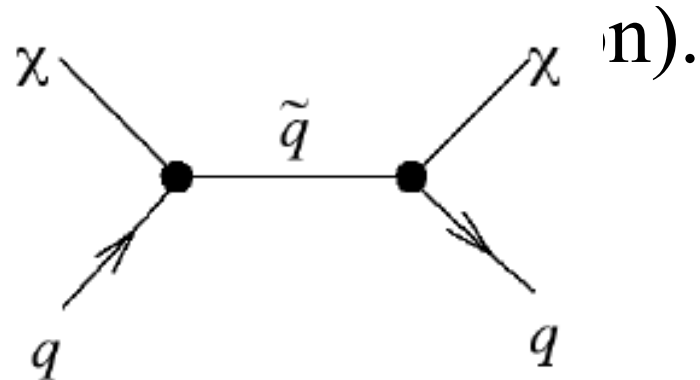
Dark Matter Searches

- Scalar elastic neutralino-nucleon scattering (DM direct detection) dominated by Higgs and squark exchange.

exc]
M()



Scalar (spin independent) couplings (tree-level)

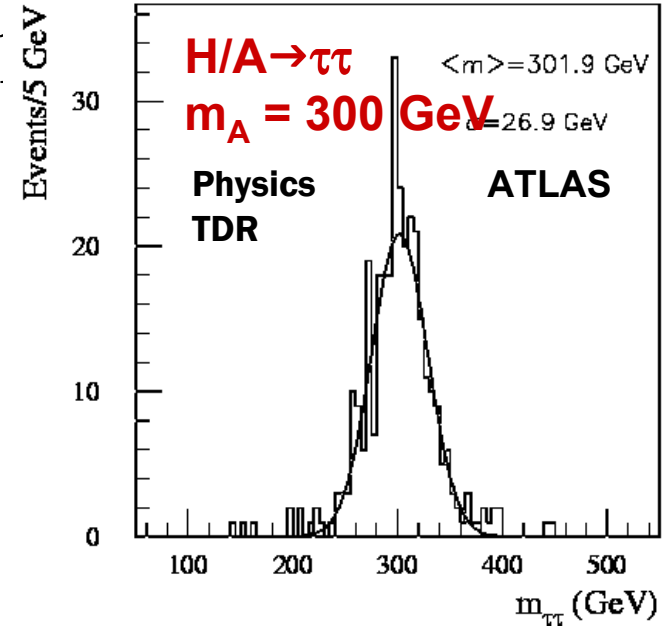
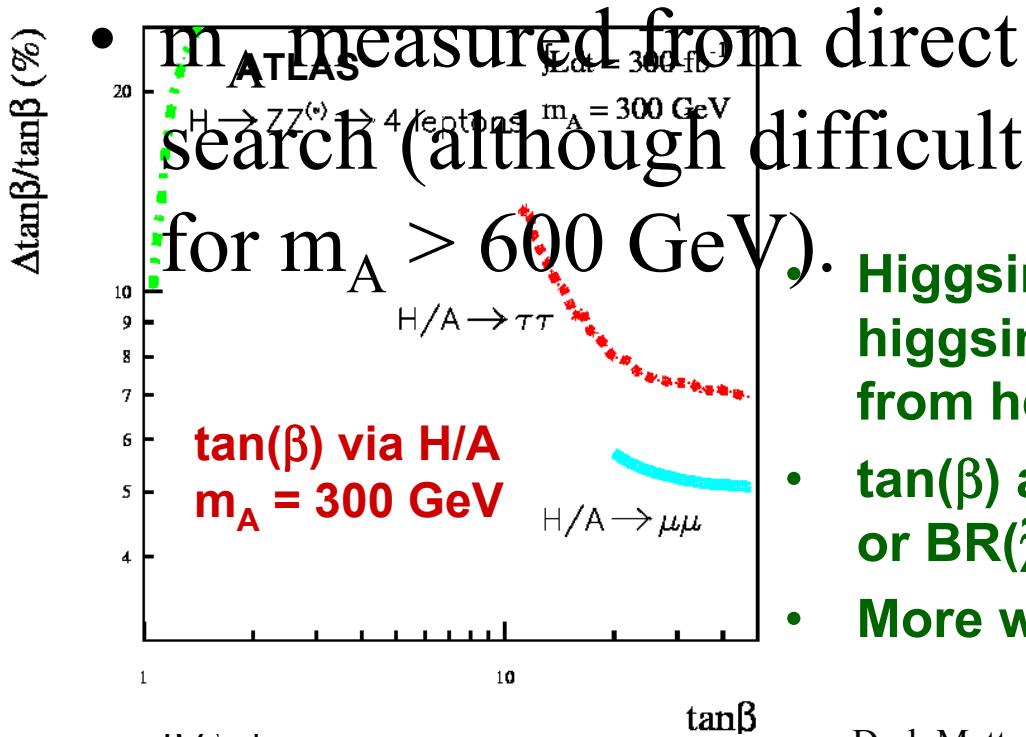


Jungman, Kamionkowski and Griest,
Phys. Rep 267:195-373 (1996)

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Other Measure

- Further input regarding the weak scale SUSY parameters needed.



- **Higgsino mass parameter μ (governs higgsino content of $\tilde{\chi}^0_1$) measurable from heavy neutralino edges.**
- **$\tan(\beta)$ accessible from $\sigma \cdot \text{BR}(H/A \rightarrow \tau\tau, \mu\mu)$ or $\text{BR}(\tilde{\chi}^0_2 \rightarrow \tau\tilde{\tau}_1) / \text{BR}(\tilde{\chi}^0_2 \rightarrow \tilde{l}l_R)$.**
- **More work needed.**

Summary

- Following a SUSY discovery ATLAS will aim to test the SUSY Dark Matter hypothesis.
- Conclusive result only possible in conjunction with astroparticle experiments (constraints on LSP life-time).
- Estimation of relic density and direct / indirect DM detection cross-sections in model-dependent scenario will be first goal. **This would be major triumph for both Particle Physics and Cosmology!**
- Less model-dependent measurements will follow.
- Ultimate goal: observation of neutralinos at

Earth from Apollo 17 (NASA)



The Fundamental Particles

Leptons

$$\begin{pmatrix} \nu_e \\ \sim 0.0 \text{ MeV } 0 e \\ 0.5 \text{ MeV } -1 e \\ e \end{pmatrix}$$

$$\begin{pmatrix} \nu_\mu \\ \sim 0.0 \text{ MeV } 0 e \\ 0.1 \text{ GeV } -1 e \\ \mu \end{pmatrix}$$

$$\begin{pmatrix} \nu_\tau \\ \sim 0.0 \text{ MeV } 0 e \\ 1.8 \text{ GeV } -1 e \\ \tau \end{pmatrix}$$

Quarks

$$\begin{pmatrix} uuu \\ 5.0 \text{ MeV } +2/3 e \\ 10.5 \text{ MeV } -1/3 e \\ ddd \end{pmatrix}$$

$$\begin{pmatrix} ccc \\ 1.3 \text{ GeV } +2/3 e \\ 0.2 \text{ GeV } -1/3 e \\ sss \end{pmatrix}$$

$$\begin{pmatrix} ttt \\ 175 \text{ GeV } +2/3 e \\ 4.3 \text{ GeV } -1/3 e \\ bbb \end{pmatrix}$$